

Chasing patterns: diffusion-driven instabilities in networked and high-order systems

Riccardo Muolo, Namur Institute for Complex Systems (naXys),
Department of Mathematics, Université de Namur (Belgium)

funded by

My background

BSc Physics and MSc Applied Mathematics
Firenze



PhD Systems Biology (1 year - quit)
Amsterdam



Teaching assistant Mathematics
Namur

Since Jan 2020
PhD Applied Mathematics (exp. 2023)
Namur



Université de Namur - naXys



Prof. Teo Carletti

Dynamics on networks and beyond group

pattern-formation, random walks,
synchronization and more



group web page



INSTITUTO DE MATEMÁTICA
Universidade Federal do Rio de Janeiro



Erasmus+

Possibility of joint Master Thesis and co-supervised PhDs

Outline

Dynamical systems on networks

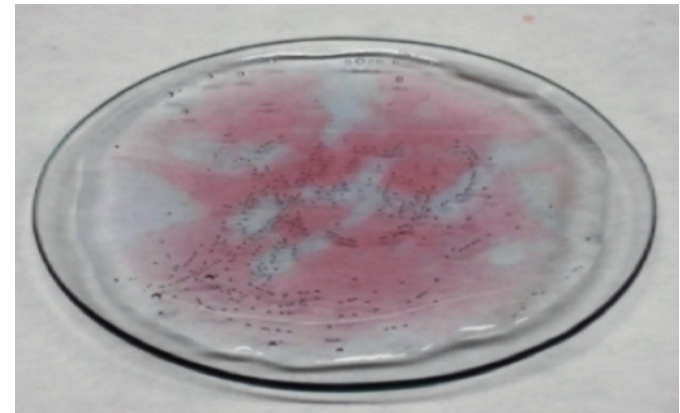
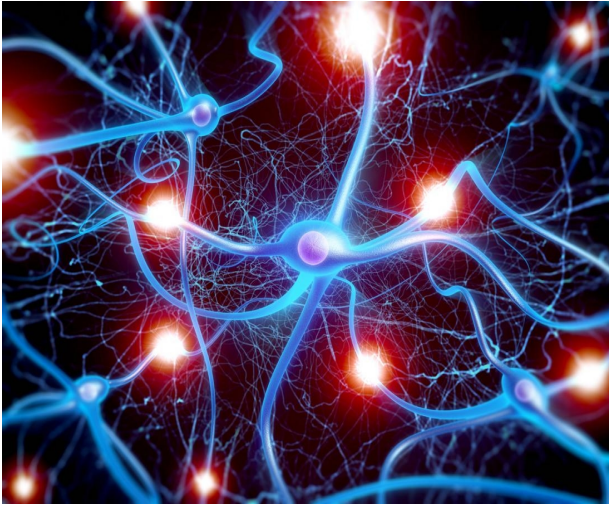
Turing theory of pattern formation

Turing theory on networks

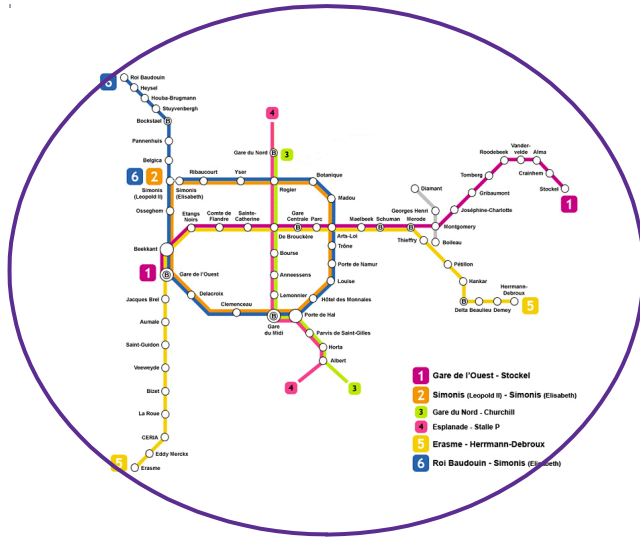
Effects of non-normality

High-order interactions

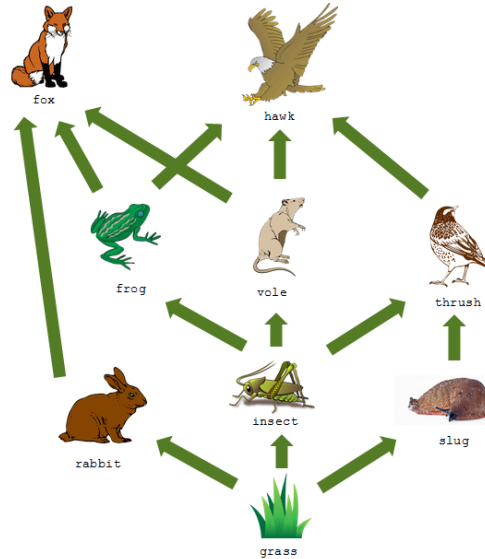
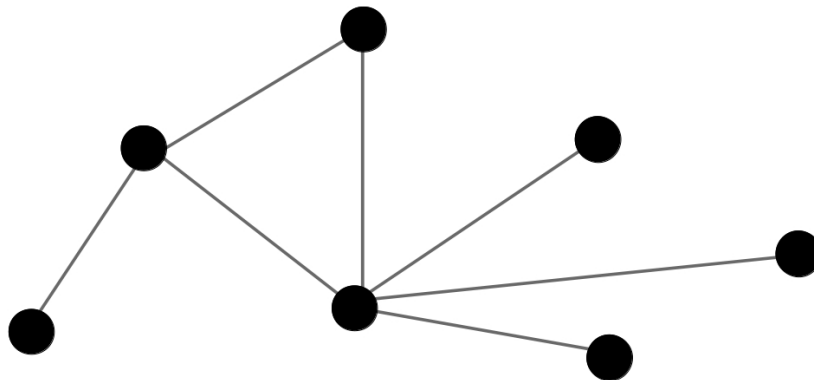
Patterns in nature



Networks



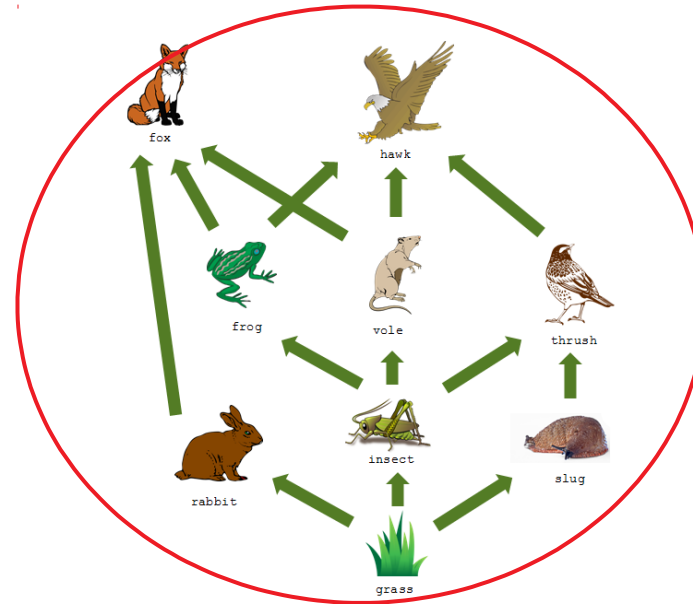
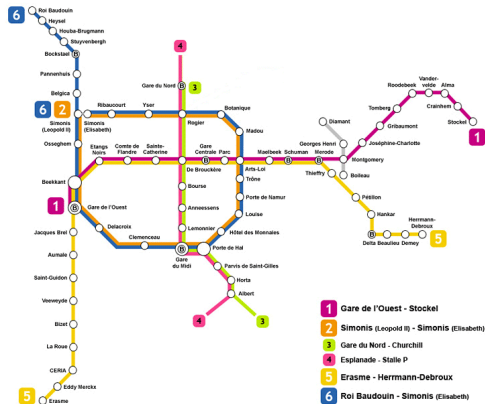
symmetric (undirected)



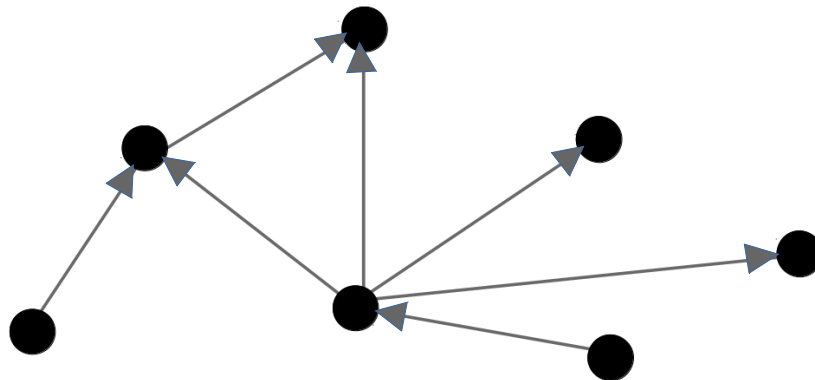
$$A_{ij} = \begin{cases} w_{ij} \in \mathbb{R}^+ \\ 0 \end{cases}$$

if there is a link
between nodes
i and j

Networks



asymmetric (directed)



$$A_{ij} = \begin{cases} w_{ij} \in \mathbb{R}^+ \\ 0 \end{cases}$$

if there is a directed link
from node j
to node i

Coupled oscillators on networks

$$\dot{\vec{x}}_i = \vec{f}(\vec{x}_i)$$

Coupled oscillators on networks

$$\dot{\vec{x}}_i = \vec{f}(\vec{x}_i)$$



dynamics of x_i



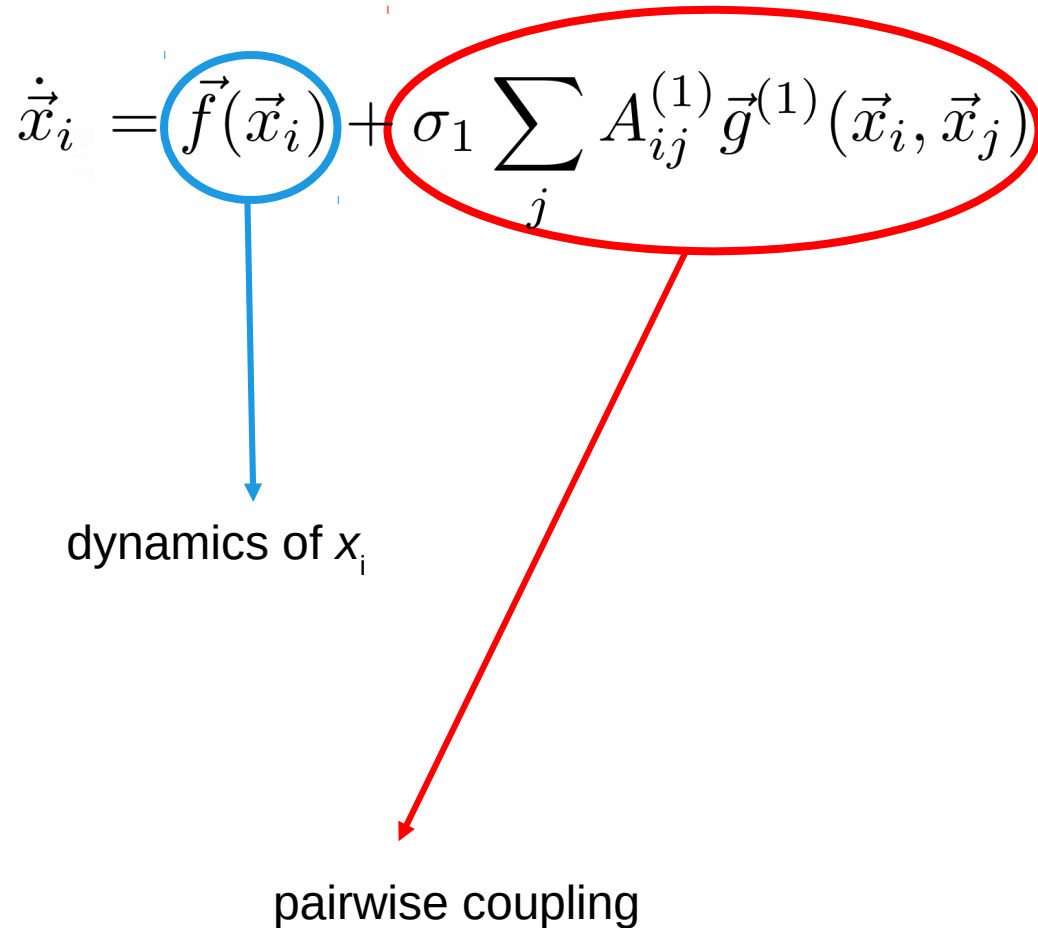
Coupled oscillators on networks

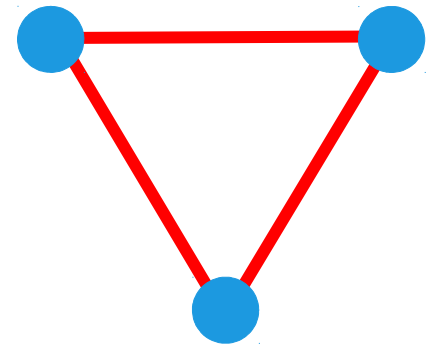
$$\dot{\vec{x}}_i = \vec{f}(\vec{x}_i) + \sigma_1 \sum_j A_{ij}^{(1)} \vec{g}^{(1)}(\vec{x}_i, \vec{x}_j)$$

dynamics of x_i



Coupled oscillators on networks

$$\dot{\vec{x}}_i = \underbrace{\vec{f}(\vec{x}_i)}_{\text{dynamics of } x_i} + \underbrace{\sigma_1 \sum_j A_{ij}^{(1)} \vec{g}^{(1)}(\vec{x}_i, \vec{x}_j)}_{\text{pairwise coupling}}$$




Master Stability Function (MSF)

VOLUME 80, NUMBER 10

PHYSICAL REVIEW LETTERS

9 MARCH 1998

Master Stability Functions for Synchronized Coupled Systems

Louis M. Pecora and Thomas L. Carroll

Code 6343, Naval Research Laboratory, Washington, D.C. 20375

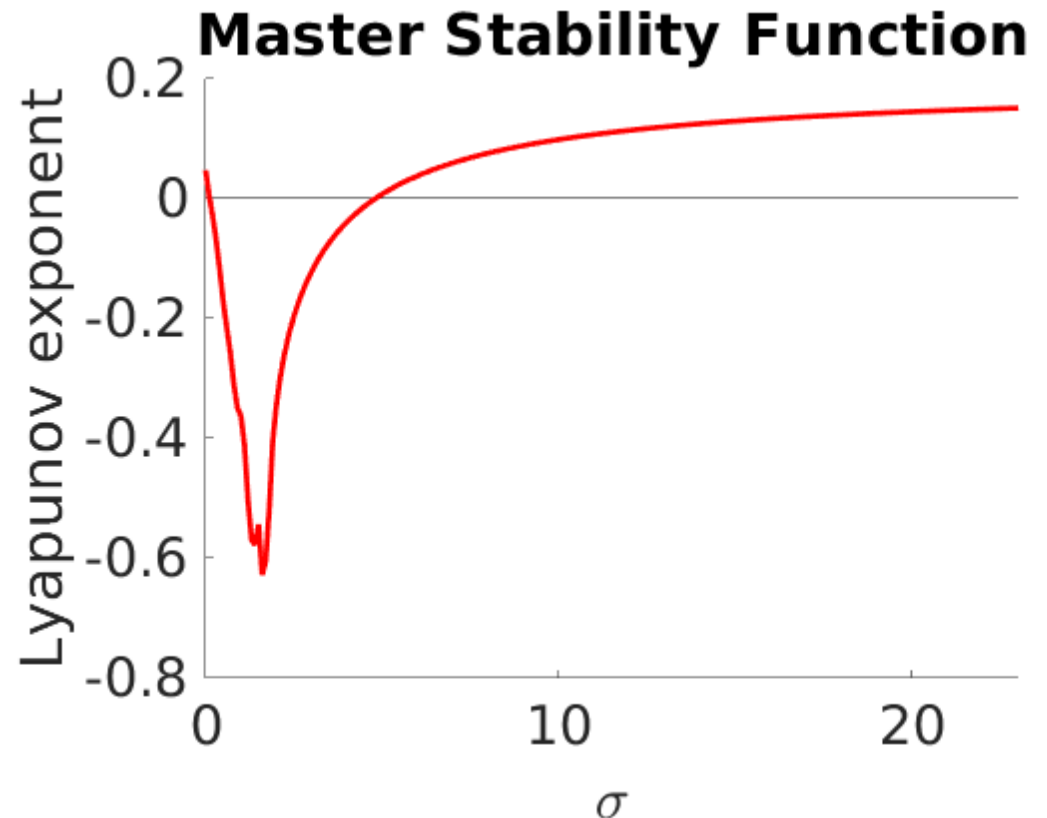
(Received 7 July 1997)

$$\dot{\xi} = [\mathbf{1}_N \otimes D\mathbf{F} + \sigma \mathbf{G} \otimes D\mathbf{H}]\xi$$

$$\dot{\xi}_k = [D\mathbf{F} + \sigma \gamma_k D\mathbf{H}]\xi_k$$

Rössler

$$\begin{cases} \dot{x}_i = -y_i - z_i + \sigma \sum_{j=1}^N L_{ij} x_j \\ \dot{y}_i = x_i + a y_i \\ \dot{z}_i = b + z_i(x_i - c) \end{cases}$$



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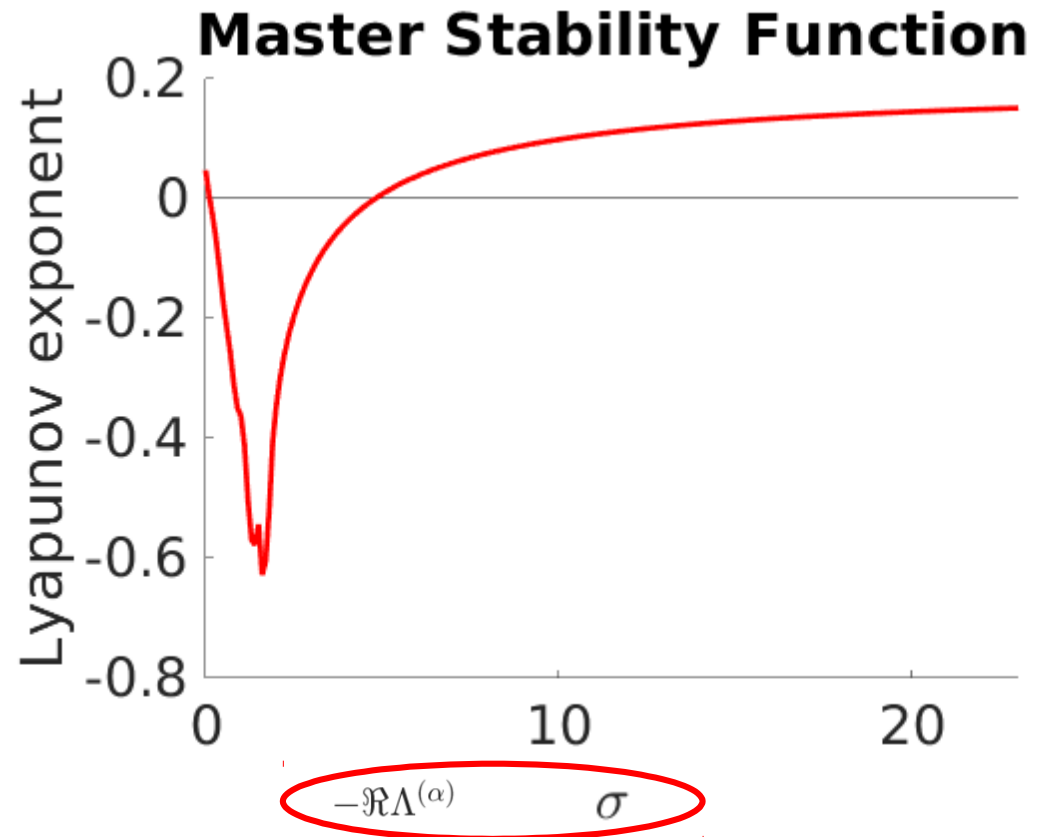
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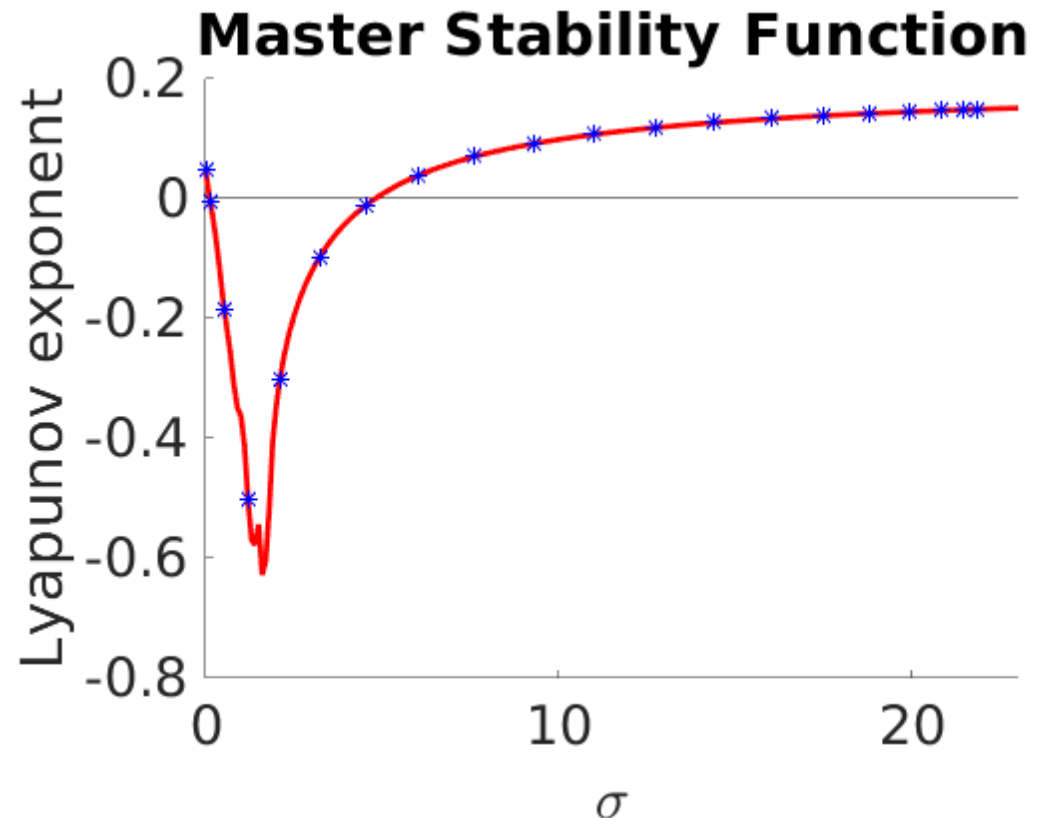
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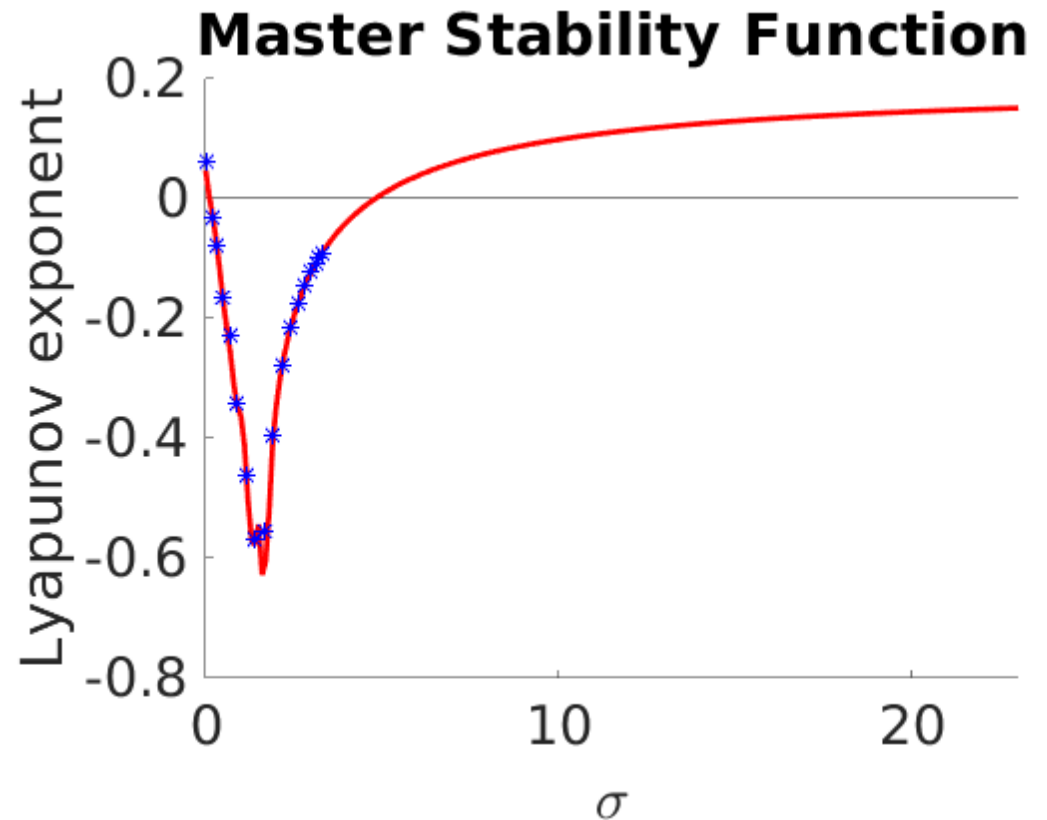
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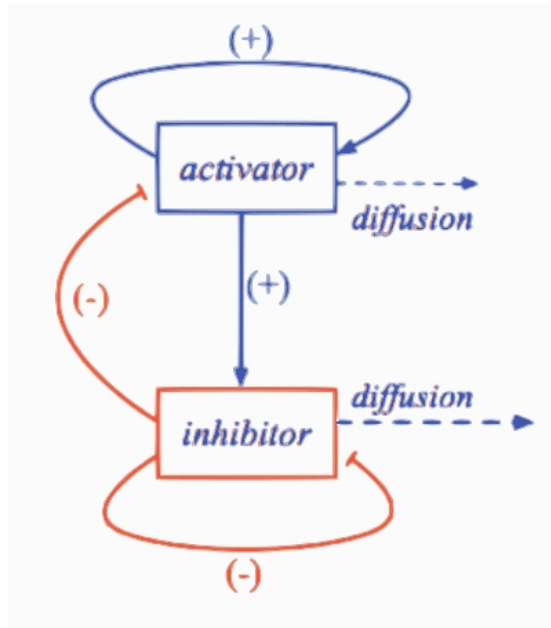
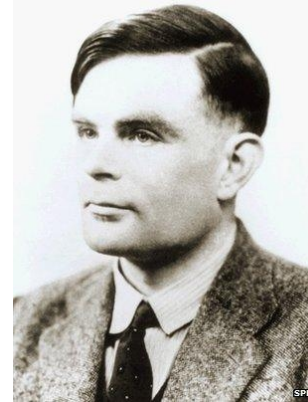


Intermezzo: Turing theory

The Chemical Basis of Morphogenesis

A. M. Turing

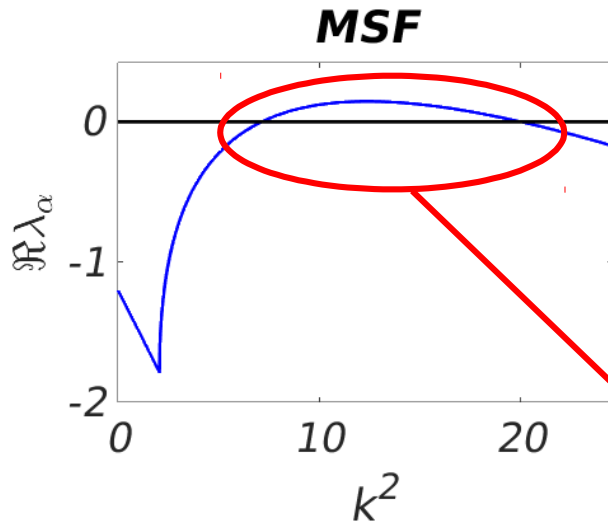
Philosophical Transactions of the Royal Society of London. Series B, Biological Sciences, Vol. 237, No. 641. (Aug. 14, 1952), pp. 37-72.



$$\begin{cases} \frac{\partial u}{\partial t}(x, t) = f(u, v) + D_u \nabla^2 u(x, t) \\ \frac{\partial v}{\partial t}(x, t) = g(u, v) + D_v \nabla^2 v(x, t) \end{cases}$$

+ boundary conditions
+ domain of the Laplacian
+

Turing theory in a nutshell



unstable modes

reaction-diffusion system of two species
(activator u and inhibitor v)

homogeneous stable state (fixed point)

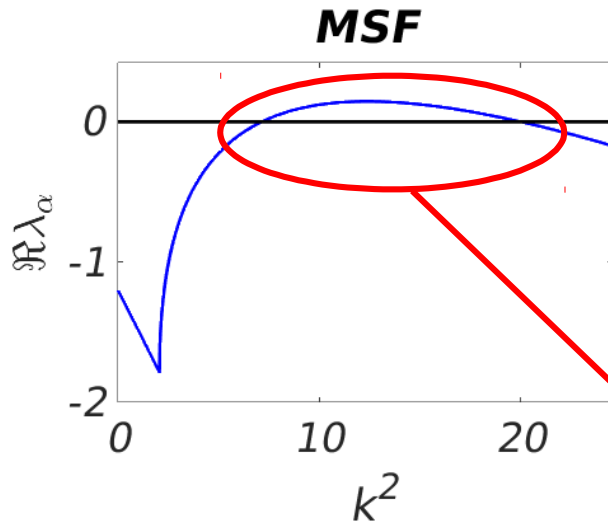
inhomogeneous perturbations

exponential instability (diffusion-driven)

patterns

$$D_v > D_u$$

Turing theory in a nutshell



unstable modes

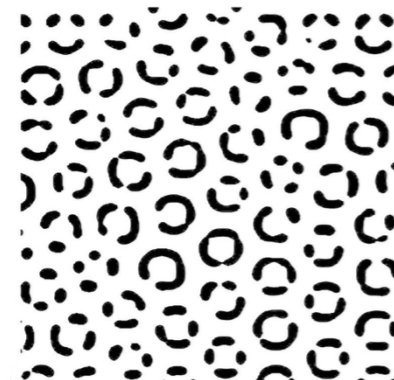
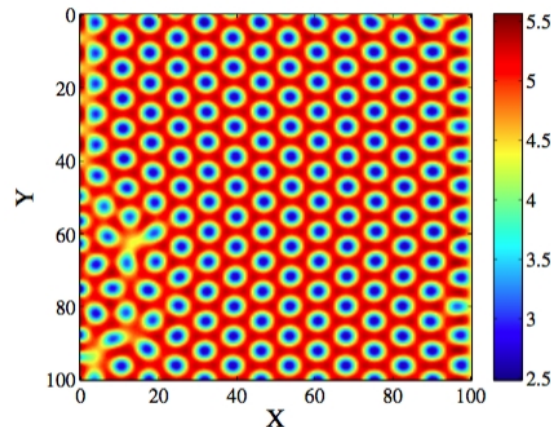
reaction-diffusion system of two species
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homogeneous stable state (fixed point)

inhomogeneous perturbations

exponential instability (diffusion-driven)

patterns $D_v > D_u$



Many applications

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RESEARCH ARTICLE

A dot-stripe Turing model of joint patterning in the tetrapod limb

Jake Cornwall Scoones¹ and Tom W. Hiscock^{2,3,*}

Article | [Published: 08 July 2021](#)

Nanoscale Turing patterns in a bismuth monolayer

[Yuki Fuseya](#) , [Hiroyasu Katsuno](#), [Kamran Behnia](#) & [Aharon Kapitulnik](#) 

[Nature Physics](#) **17**, 1031–1036 (2021) | [Cite this article](#)

PHYSICAL REVIEW RESEARCH **3**, 023241 (2021)

scientific reports

OPEN

Turing instability in quantum activator–inhibitor systems

Yuzuru Kato  & Hiroya Nakao 

Linguistic evolution driven by network heterogeneity and the Turing mechanism

Sayat Mimar ¹, Mariamo Mussa Juane², Jorge Mira ³, Juyong Park⁴, Alberto P. Muñozuri ² and Gourab Ghoshal ^{1,*}

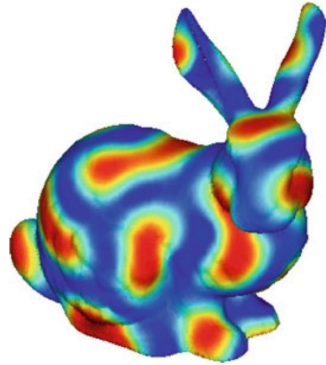
¹Department of Physics & Astronomy, University of Rochester, Rochester, New York 14607, USA

²Group of Nonlinear Physics, Universidade de Santiago de Compostela, 15782 Santiago de Compostela, Spain

³Departamento de Física Aplicada, Universidade de Santiago de Compostela, 15782 Santiago de Compostela, Spain

⁴Graduate School of Culture Technology, Korea Advanced Institute of Science and Technology, Daejeon, 305-701, Korea

Many applications



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scientific reports

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⁴Graduate School of Culture Technology, Korea Advanced Institute of Science and Technology, Daejeon, 305-701, Korea

Extension on networks

J. theor. Biol. (1971) **32**, 507–537

ARTICLES

PUBLISHED ONLINE: 25 APRIL 2010 | DOI: 10.1038/NPHYS1651

nature
physics

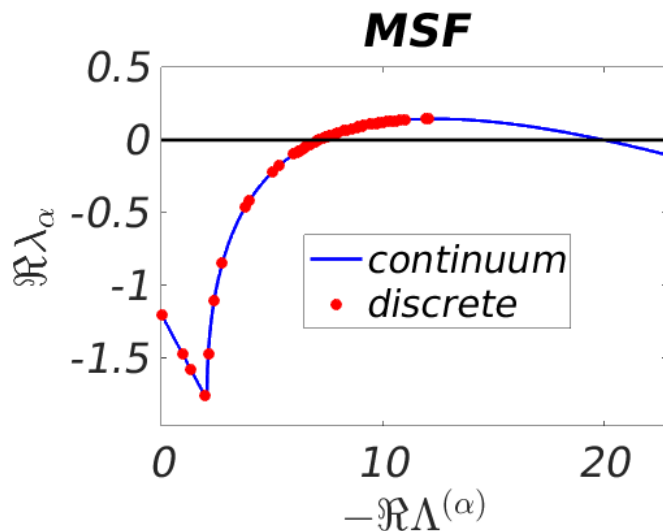
Turing patterns in network-organized activator-inhibitor systems

Hiroya Nakao^{1,2*} and Alexander S. Mikhailov^{3*}

Instability and Dynamic Pattern in Cellular Networks

H. G. OTHMER[†] AND L. E. SCRIVEN[‡]

Department of Chemical Engineering and Materials Science
Institute of Technology, University of Minnesota,
Minneapolis, Minnesota 55455, U.S.A.



$$L_{ij} = A_{ij} - k_i \delta_{ij}$$

$$\begin{cases} \frac{du_i}{dt} = f(u_i, v_i) + D_u \sum_{j=1}^{\text{nodes}} L_{ij} u_j \\ \frac{dv_i}{dt} = g(u_i, v_i) + D_v \sum_{j=1}^{\text{nodes}} L_{ij} v_j \end{cases}$$

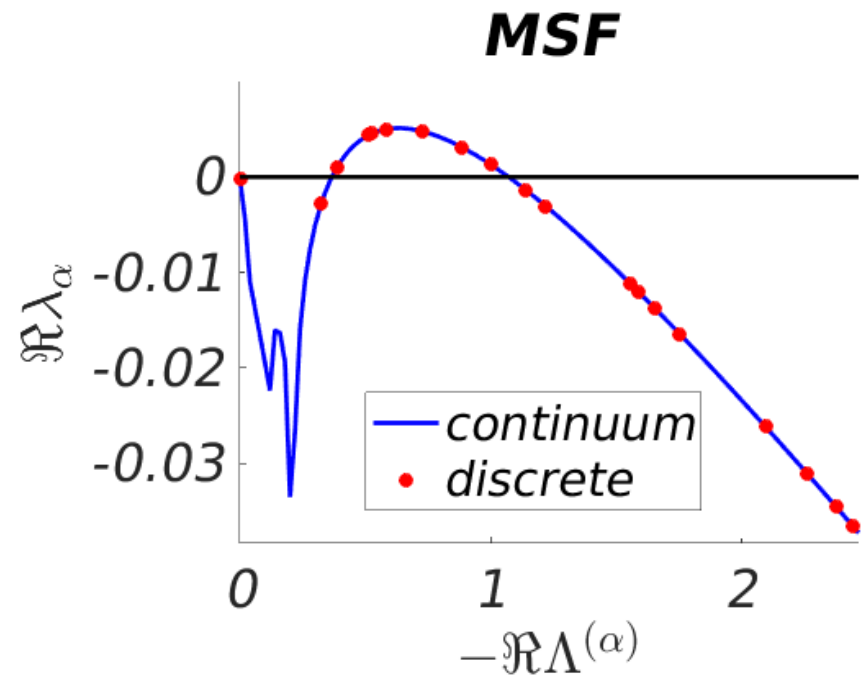
Turing-like instability

PHYSICAL REVIEW E **92**, 022818 (2015)

Turing-like instabilities from a limit cycle

Joseph D. Challenger,^{1,2} Raffaella Burioni,³ and Duccio Fanelli²

the homogeneous stable
state is a **limit cycle**

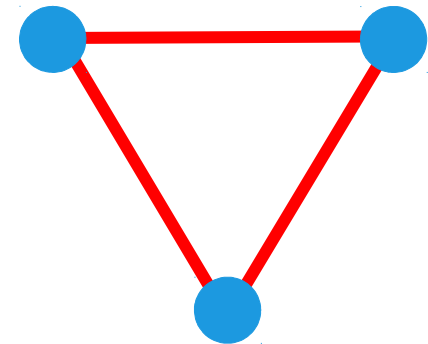


Coupled oscillators on networks (!)

$$\dot{\vec{x}}_i = \underbrace{\vec{f}(\vec{x}_i)}_{\text{blue circle}} + \underbrace{\sigma_1 \sum_j A_{ij}^{(1)} \vec{g}^{(1)}(\vec{x}_i, \vec{x}_j)}_{\text{red oval}}$$

↓
dynamics of x_i
chaotic, oscillating,
Fixed point, etc...

↘
pairwise coupling
in general diffusive like

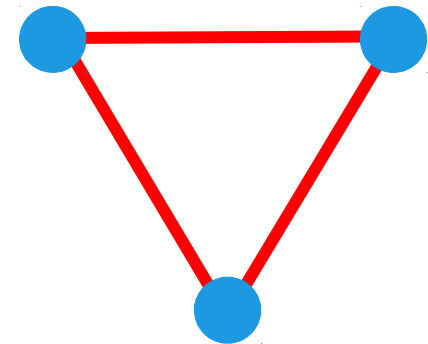
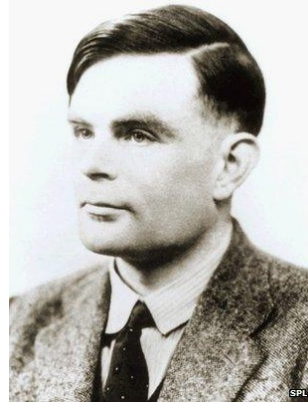


Coupled oscillators on networks (!)

$$\dot{\vec{x}}_i = \underbrace{\vec{f}(\vec{x}_i)}_{\text{dynamics of } x_i} + \underbrace{\sigma_1 \sum_j A_{ij}^{(1)} \vec{g}^{(1)}(\vec{x}_i, \vec{x}_j)}_{\text{pairwise coupling in general diffusive like}}$$

dynamics of x_i
chaotic, oscillating,
Fixed point, etc...

pairwise coupling
in general diffusive like



Extensions of Turing theory

PHYSICAL REVIEW E **92**, 022818 (2015)

Turing-like instabilities from a limit cycle

Joseph D. Challenger,^{1,2} Raffaella Burioni,³ and Duccio Fanelli²



ARTICLE

Received 5 Feb 2014 | Accepted 26 Jun 2014 | Published 31 Jul 2014

DOI: 10.1038/ncomms5517

The theory of pattern formation on directed networks

Malbor Asllani^{1,2}, Joseph D. Challenger², Francesco Saverio Pavone^{2,3,4}, Leonardo Sacconi^{3,4} & Duccio Fanelli²

Journal of Theoretical Biology 480 (2019) 81–91



Contents lists available at ScienceDirect

Journal of Theoretical Biology

journal homepage: www.elsevier.com/locate/jtb



Patterns of non-normality in networked systems

Riccardo Muolo^a, Malbor Asllani^{b,c,*}, Duccio Fanelli^d, Philip K. Maini^b, Timoteo Carletti^e



Journal of Physics: Complexity

OPEN ACCESS



PAPER

Finite propagation enhances Turing patterns in reaction–diffusion networked systems

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27 August 2021

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29 September 2021

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* Author to whom any correspondence should be addressed.

E-mail: timoteo.carletti@unamur.be

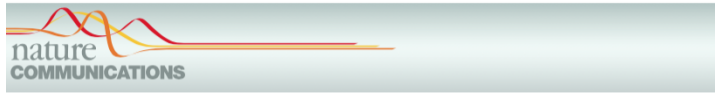
Chaos, Solitons and Fractals 134 (2020) 109707

Generalized patterns from local and non local reactions

Giulia Cencetti^a, Federico Battiston^b, Timoteo Carletti^c, Duccio Fanelli^{d,*}



Directed Networks



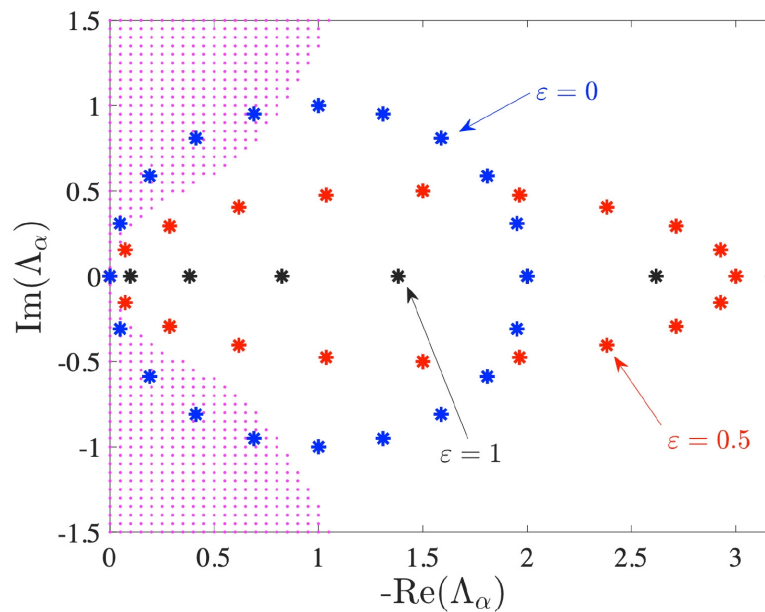
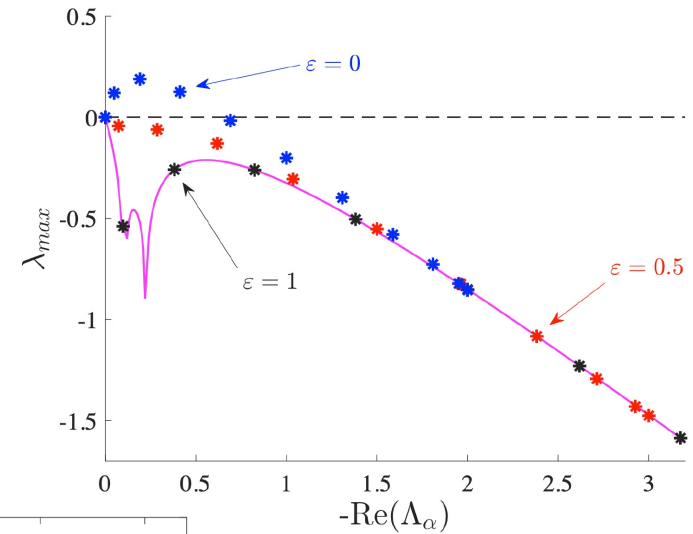
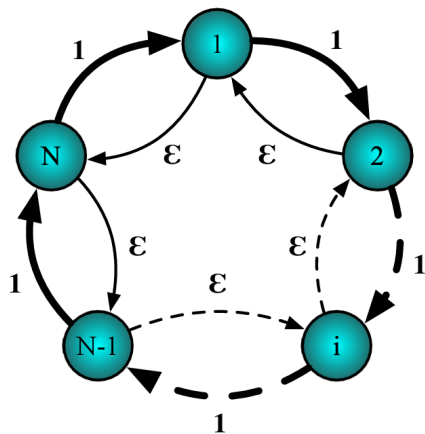
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The theory of pattern formation on directed networks

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Non-normal Networks

A network is non-normal
when $A^*A \neq AA^*$

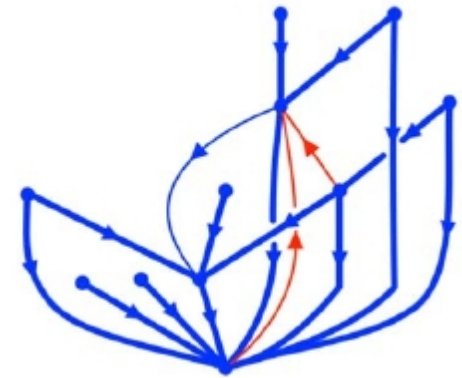
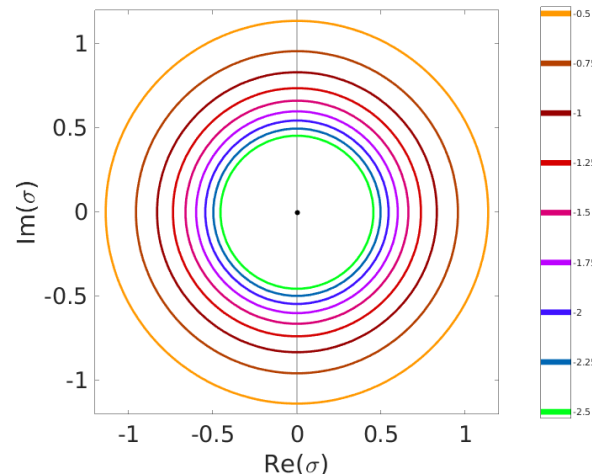
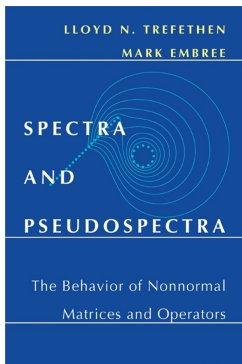
RESEARCH ARTICLE | NETWORK SCIENCE

Structure and dynamical behavior of non-normal networks

 Malbor Asllani^{1,2},  Renaud Lambiotte¹ and  Timoteo Carletti^{2,*}

+ See all authors and affiliations

Science Advances 12 Dec 2018;
Vol. 4, no. 12, eaau9403
DOI: 10.1126/sciadv.aau9403

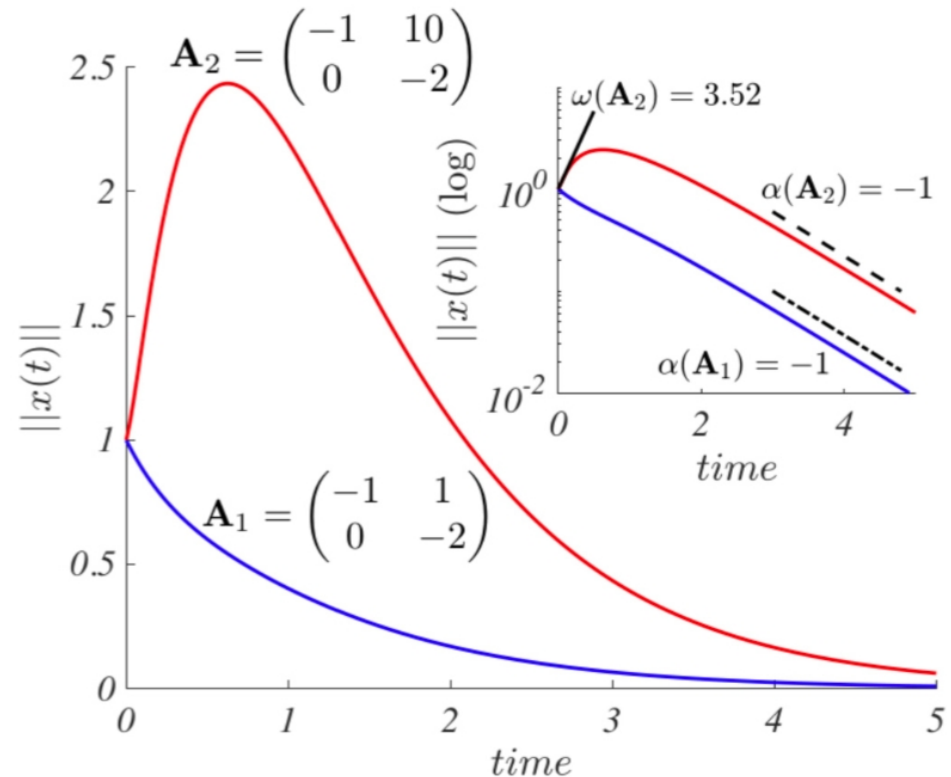


$$\beta\text{-pseudospectrum } \lambda_{\beta}(A) \rightarrow \lambda(A+B) \text{ } ||B|| < \beta$$

Effects on the dynamics

ω largest eigenvalue of the
hermitian (symmetric) part

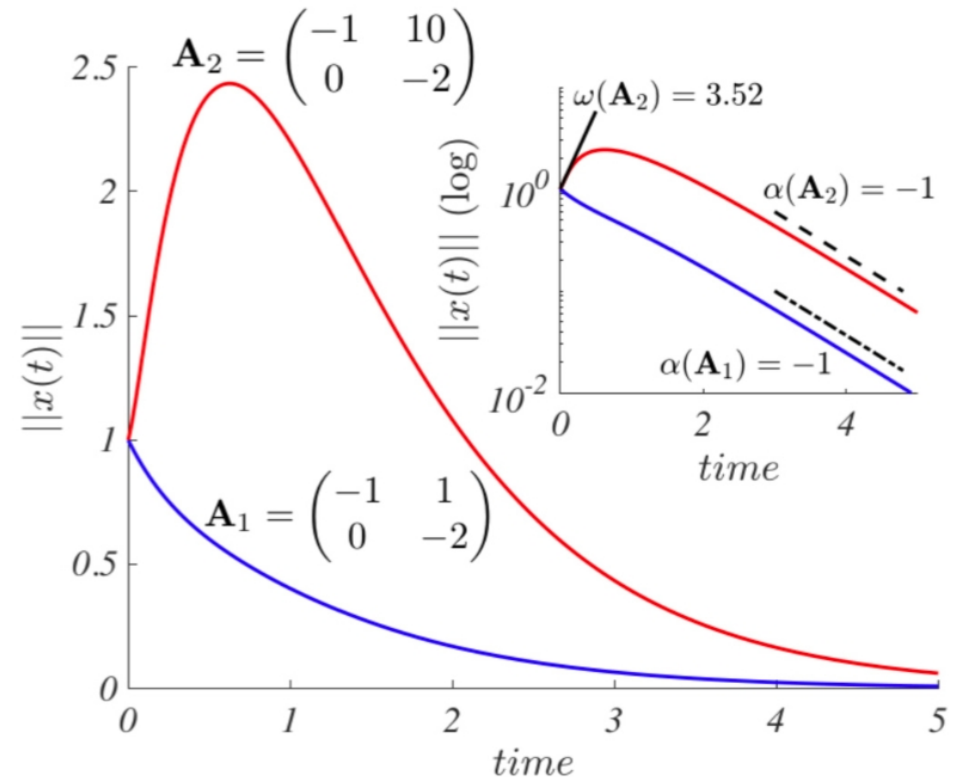
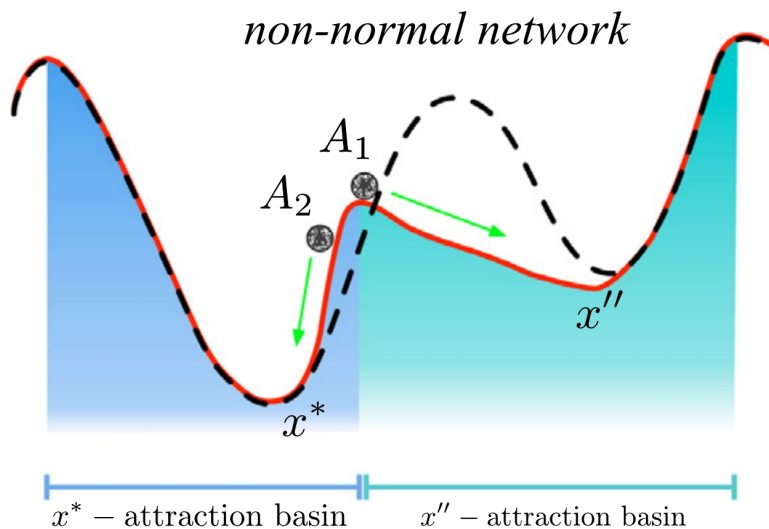
Linear dynamics



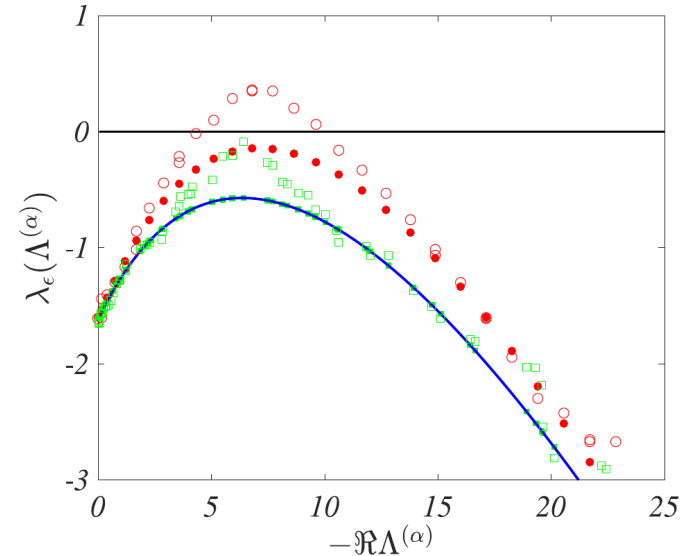
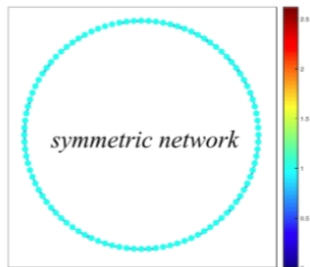
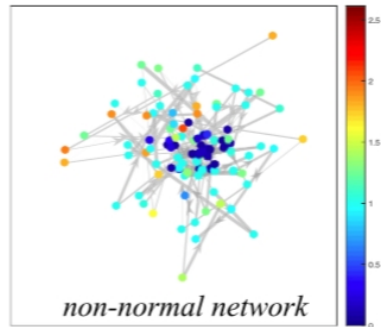
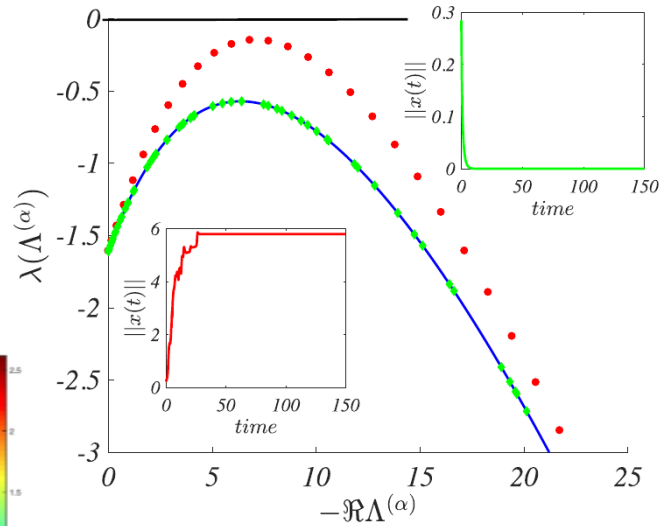
Effects on the dynamics

Linear dynamics

Non-linear dynamics



Patterns of Non-normality



Journal of Theoretical Biology 480 (2019) 81–91



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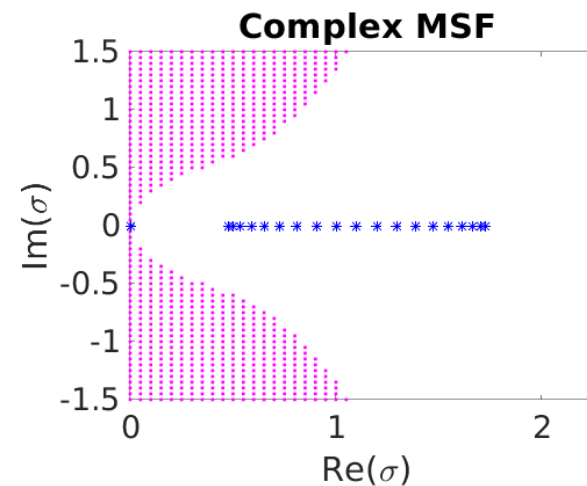
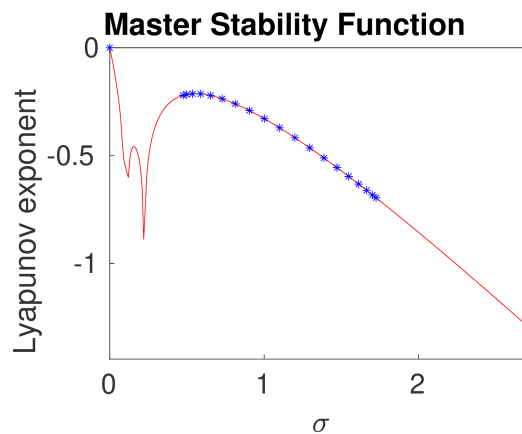
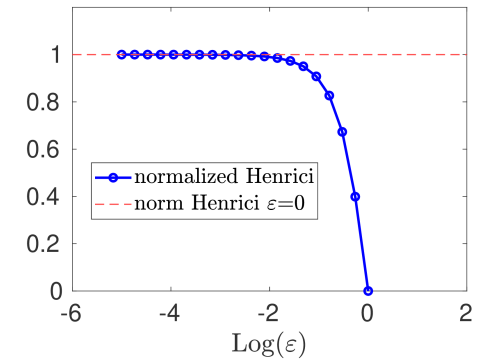
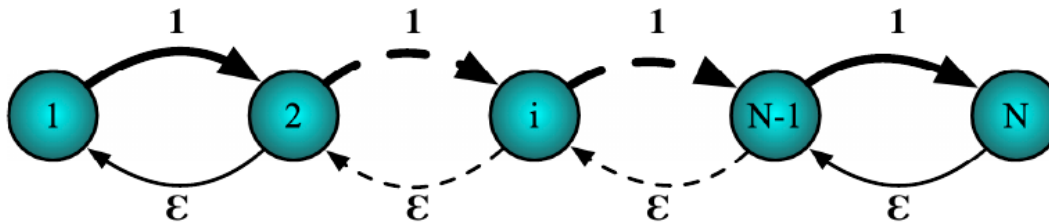


Synchronization and non-normality

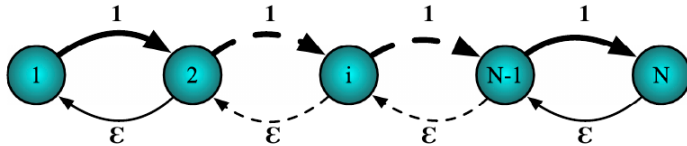
Article

Synchronization Dynamics in Non-Normal Networks: The Trade-Off for Optimality

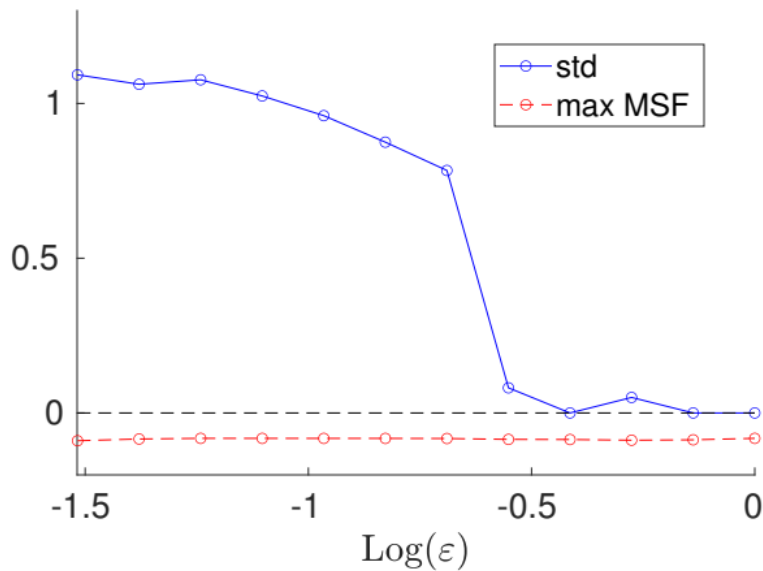
Riccardo Muolo ^{1,*} , Timoteo Carletti ¹ , James P. Gleeson ²  and Malbor Asllani ^{1,2} 



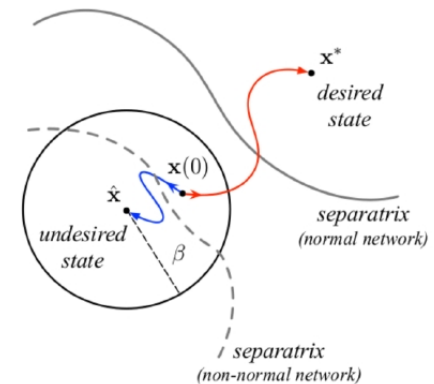
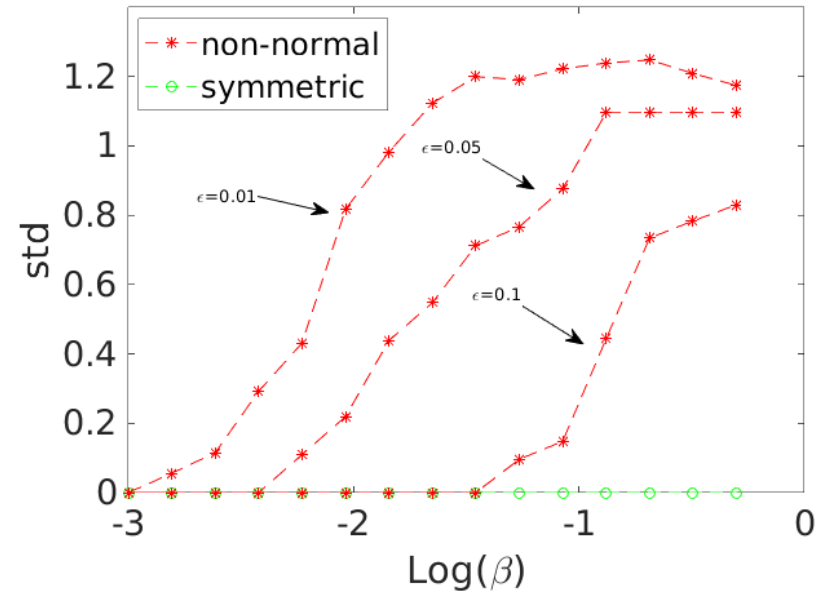
De-synchronization



MSF vs non-linear behavior



basin of attraction



Debate on optimal networks



Article

Synchronization Dynamics in Non-Normal Networks: The Trade-Off for Optimality

Riccardo Muolo ^{1,*}, Timoteo Carletti ¹, James P. Gleeson ² and Malbor Asllani ^{1,2}

Comment on “Synchronization dynamics in non-normal networks: the trade-off for optimality”

Takashi Nishikawa,^{1,2,*} Adilson E. Motter,^{1,2} and Louis M. Pecora³

Reply to Comment on “Synchronization dynamics in non-normal networks: the trade-off for optimality”

Riccardo Muolo¹, Timoteo Carletti¹, James P. Gleeson², Malbor Asllani³

Comment on “Synchronization dynamics in non-normal networks: the trade-off for optimality”

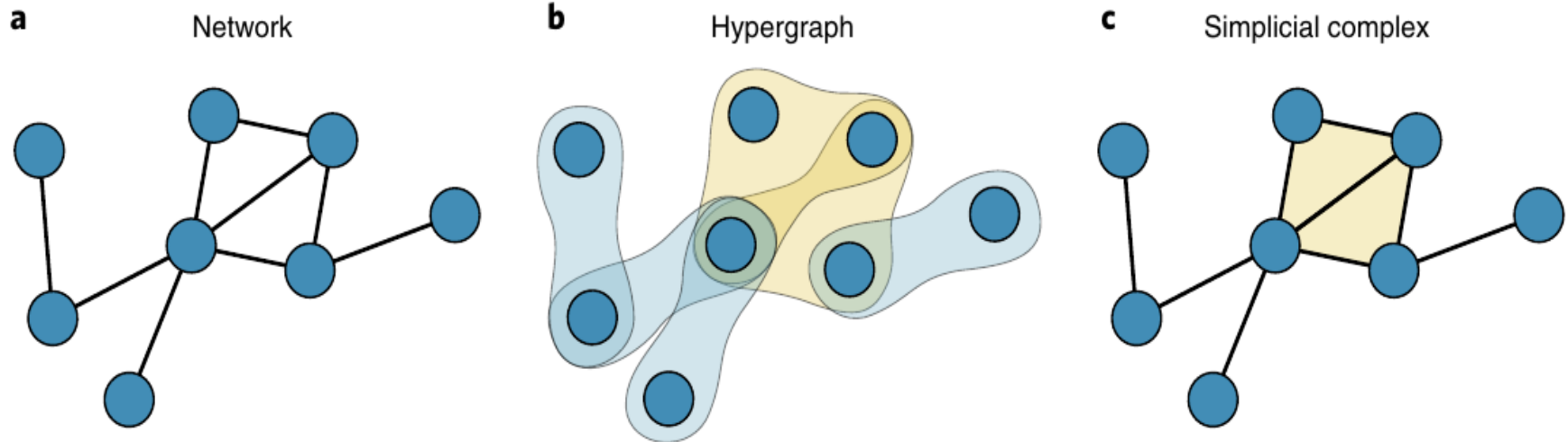
Francesco Sorrentino and Chad Nathe

Non-normality, optimality and synchronization

Jeremie Fish*

Erik M. Bollt

High-order Structures

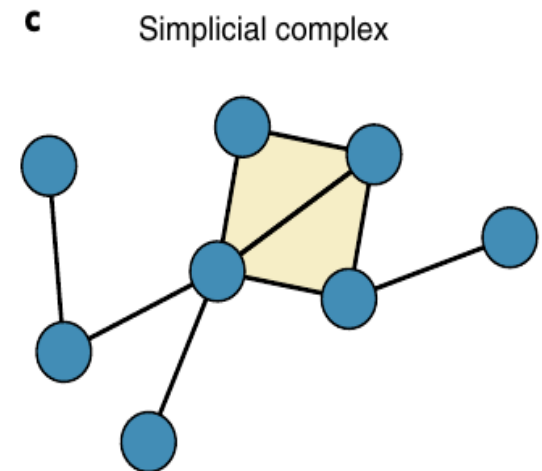
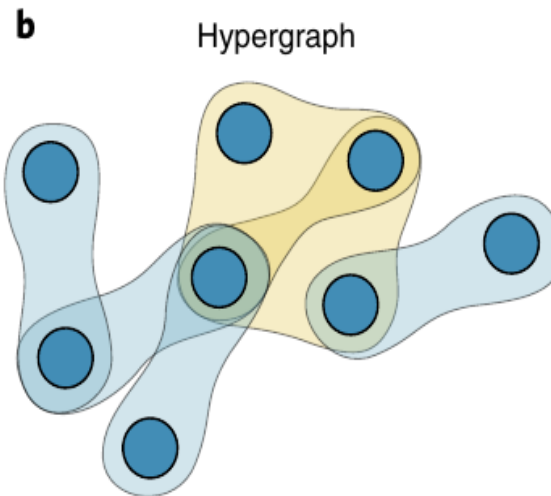
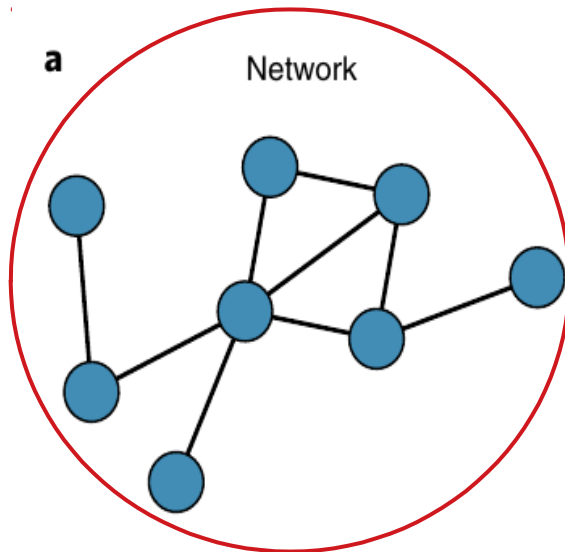


Battiston et al., Nat. Phys., 2021

High-order Structures

adjacency matrix

adjacency tensors

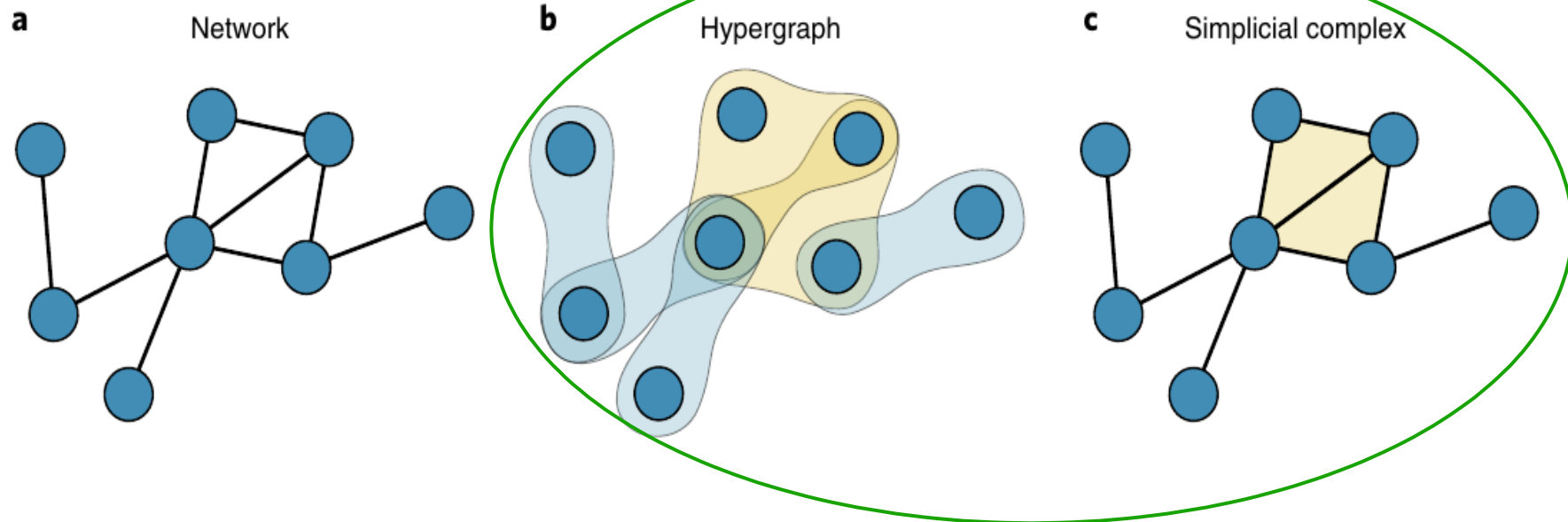


Battiston et al., Nat. Phys., 2021

High-order Structures

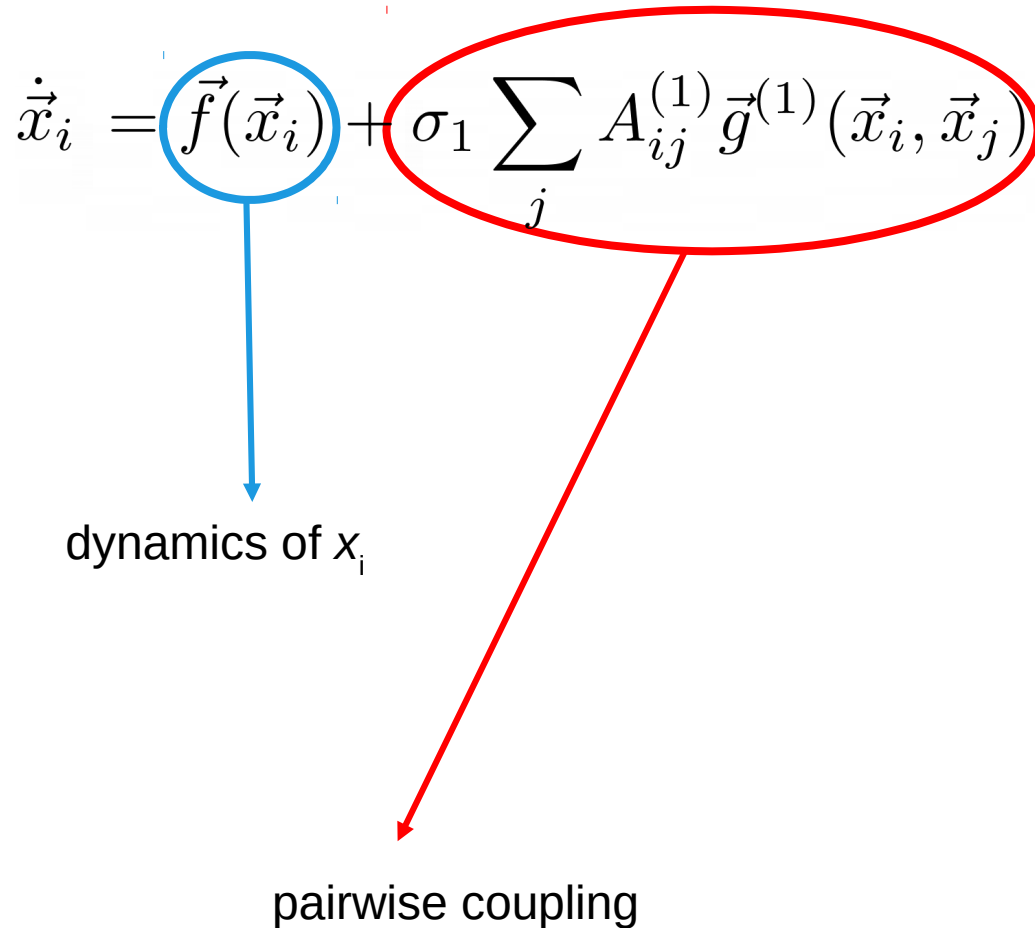
adjacency matrix

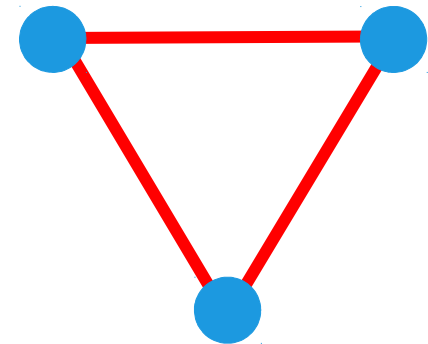
adjacency tensors



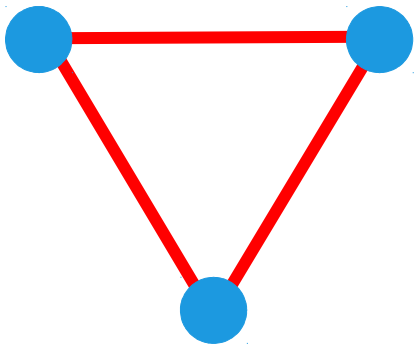
Battiston et al., Nat. Phys., 2021

High-order coupled oscillators

$$\dot{\vec{x}}_i = \underbrace{\vec{f}(\vec{x}_i)}_{\text{dynamics of } x_i} + \underbrace{\sigma_1 \sum_j A_{ij}^{(1)} \vec{g}^{(1)}(\vec{x}_i, \vec{x}_j)}_{\text{pairwise coupling}}$$


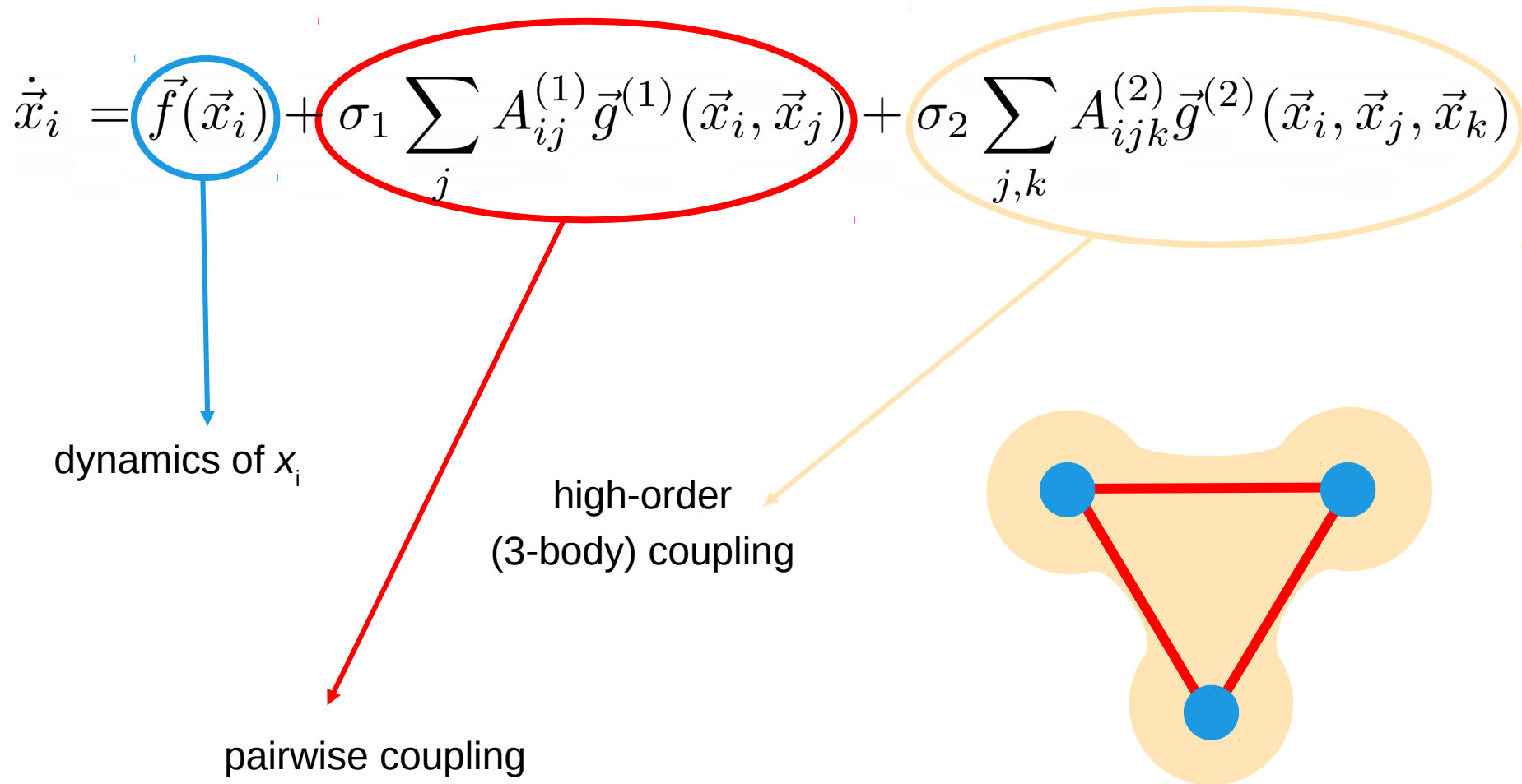


High-order coupled oscillators

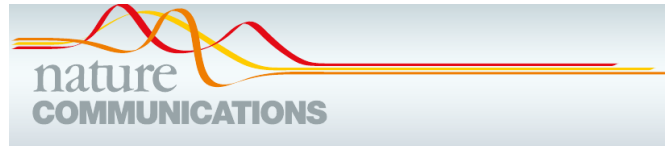
$$\dot{\vec{x}}_i = \underbrace{\vec{f}(\vec{x}_i)}_{\text{dynamics of } x_i} + \underbrace{\sigma_1 \sum_j A_{ij}^{(1)} \vec{g}^{(1)}(\vec{x}_i, \vec{x}_j)}_{\text{pairwise coupling}} + \sigma_2 \sum_{j,k} A_{ijk}^{(2)} \vec{g}^{(2)}(\vec{x}_i, \vec{x}_j, \vec{x}_k)$$


The diagram illustrates the components of the equation for high-order coupled oscillators. The first term, $\vec{f}(\vec{x}_i)$, is circled in blue and labeled "dynamics of x_i ". The second term, $\sigma_1 \sum_j A_{ij}^{(1)} \vec{g}^{(1)}(\vec{x}_i, \vec{x}_j)$, is circled in red and labeled "pairwise coupling". A red arrow points from this term to a diagram of a triangular network of three blue nodes connected by red lines, representing pairwise coupling. The third term, $\sigma_2 \sum_{j,k} A_{ijk}^{(2)} \vec{g}^{(2)}(\vec{x}_i, \vec{x}_j, \vec{x}_k)$, represents higher-order coupling.

High-order coupled oscillators



Dynamics on high-order structures



ARTICLE



<https://doi.org/10.1038/s41467-021-21486-9>

OPEN

Stability of synchronization in simplicial complexes

L. V. Gambuzza^{1,12}, F. Di Patti^{2,12}, L. Gallo^{3,4,12}, S. Lepri², M. Romance⁵, R. Criado⁵, M. Frasca^{1,6,13}, V. Latora^{3,4,7,8,13} & S. Boccaletti^{2,9,10,11,13}

J.Phys.Complex. **1** (2020) 035006 (16pp)

Journal of Physics: Complexity

PAPER

Dynamical systems on hypergraphs

Timoteo Carletti^{1,4}, Duccio Fanelli² and Sara Nicoletti^{2,3}

Turing theory in high-order systems

arXiv > nlin > arXiv:2207.03985

Nonlinear Sciences > Pattern Formation and Solitons

[Submitted on 8 Jul 2022]

Turing patterns in systems with high-order interactions

Riccardo Muolo, Luca Gallo, Vito Latora, Mattia Frasca, Timoteo Carletti

$$\begin{aligned}\dot{u}_i &= f_1(u_i, v_i) + \sigma_1 D_u^{(1)} \sum_{j_1=1}^N A_{ij_1}^{(1)} (h_1^{(1)}(u_{j_1}) - h_1^{(1)}(u_i)) \\ &\quad + \sigma_2 D_u^{(2)} \sum_{j_1=1}^N \sum_{j_2=1}^N A_{ij_1 j_2}^{(2)} (h_1^{(2)}(u_{j_1}, u_{j_2}) - h_1^{(2)}(u_i, u_i)) \\ \dot{v}_i &= f_2(u_i, v_i) + \sigma_1 D_v^{(1)} \sum_{j_1=1}^N A_{ij_1}^{(1)} (h_2^{(1)}(v_{j_1}) - h_2^{(1)}(v_i)) \\ &\quad + \sigma_2 D_v^{(2)} \sum_{j_1=1}^N \sum_{j_2=1}^N A_{ij_1 j_2}^{(2)} (h_2^{(2)}(v_{j_1}, v_{j_2}) - h_2^{(2)}(v_i, v_i))\end{aligned}$$

Natural coupling

$$\frac{d}{dt}\vec{\xi} = \left(\mathbb{I}_N \otimes \mathbf{J}_0 + \sigma_1 \mathbf{L}^{(1)} \otimes \mathbf{J}_{H^{(1)}} + \sigma_2 \mathbf{L}^{(2)} \otimes \mathbf{J}_{H^{(2)}} \right) \vec{\xi}$$

Natural coupling

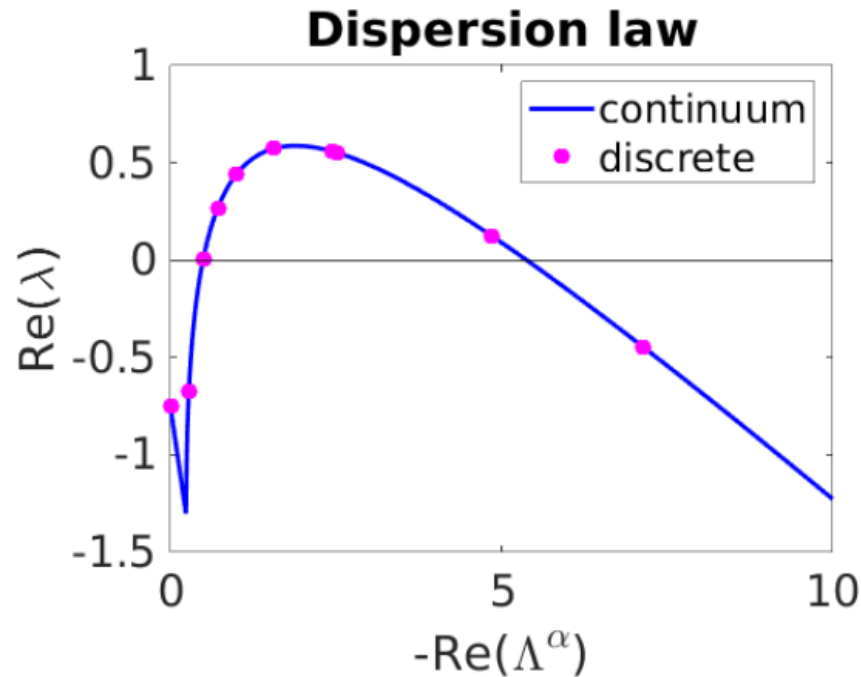
$$\frac{d}{dt}\vec{\xi} = \left(\mathbb{I}_N \otimes \mathbf{J}_0 + \sigma_1 \mathbf{L}^{(1)} \otimes \mathbf{J}_{H^{(1)}} + \sigma_2 \mathbf{L}^{(2)} \otimes \mathbf{J}_{H^{(2)}} \right) \vec{\xi}$$

$$\vec{h}^{(d)}(\vec{x}, \dots, \vec{x}) = \dots = \vec{h}^{(2)}(\vec{x}, \vec{x}) = \vec{h}^{(1)}(\vec{x})$$

same diffusion coefficients
for every order

$$\frac{d}{dt}\vec{\xi} = \left[\mathbb{I}_N \otimes \mathbf{J}_0 + \left(\sigma_1 \mathbf{L}^{(1)} + \sigma_2 \mathbf{L}^{(2)} \right) \otimes \mathbf{J}_{H^{(1)}} \right] \vec{\xi}$$

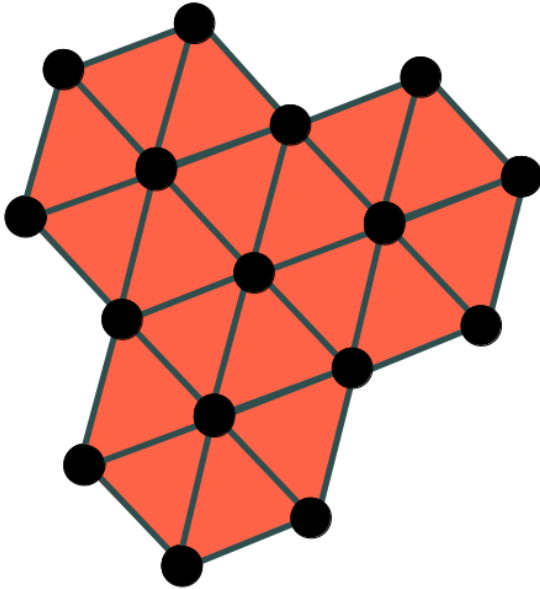
Natural coupling



same diffusion coefficients
for every order

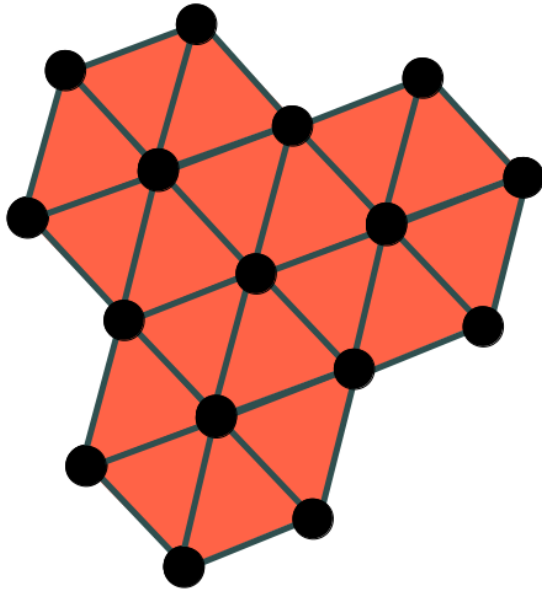
$$\frac{d}{dt}\vec{\xi} = \left[\mathbb{I}_N \otimes \mathbf{J}_0 + \left(\sigma_1 \mathbf{L}^{(1)} + \sigma_2 \mathbf{L}^{(2)} \right) \otimes \mathbf{J}_{H^{(1)}} \right] \vec{\xi}$$

Regular topologies

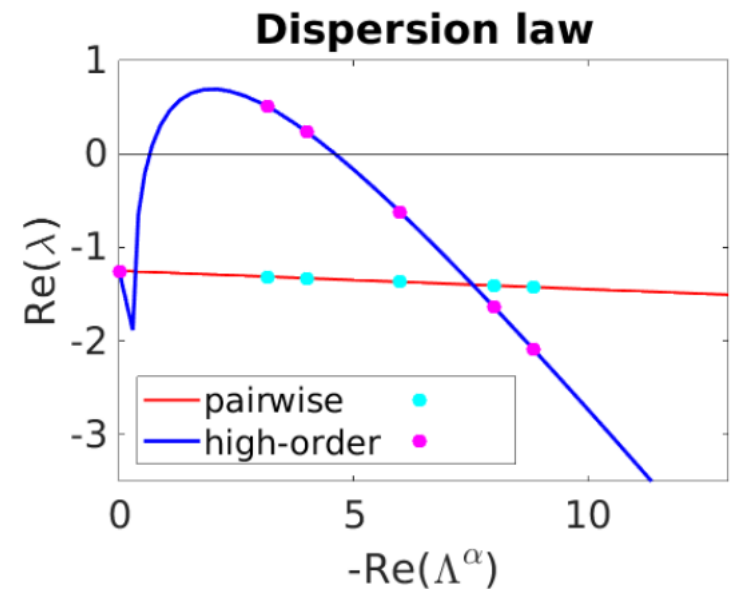


$$\mathbf{L}^{(2)} = 2\mathbf{L}^{(1)}$$

Regular topologies



$$\mathbf{L}^{(2)} = 2\mathbf{L}^{(1)}$$

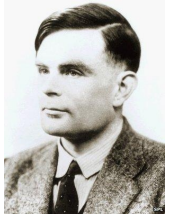


$$\frac{d}{dt}\vec{\xi} = \left(\mathbb{I}_N \otimes \mathbf{J}_0 + \mathbf{L}^{(1)} \otimes (\sigma_1 \mathbf{J}_{H^{(1)}} + 2\sigma_2 \mathbf{J}_{H^{(2)}}) \right) \vec{\xi}$$

Thank you for your attention

Questions?

Take Home Messages



synchronization of coupled oscillators can be studied in the framework of Turing pattern formation (diffusive coupling)

the topology greatly affects the stability of the synchronized state and a linear stability analysis may fail in its predictions

the theory can be extended in the framework of high-order interactions



LinkedIn Riccardo Muolo

