#### On a Dyson fractional Brownian motion

Victor Pérez-Abreu Guanajuato, Mexico

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#### Contents

- 0. Recalling pioneering work on random matrices
  - a. Wigner's law for Gaussian ensembles
  - b. Marchenko-Pastur theorem for covariance matrices
- I. From Matrix Brownian motion to Free Brownian motion
  - A. Hermitian Brownian motion
  - B. Dyson Brownian Motion: **SDE for non-colliding eigenvalues**
  - D. Spectral empirical processes and the functional Wigner's law
  - E. Law of Free Brownian motion: limiting measure-valued family
- II. Free Brownian motion and the non-commutatitve fractional Brownian motion
- III. Non-Colliding eigenvalues and SDE for other matrix processes:
  - A. Fractional matrix Brownian motion and the non-commutatitve fractional Brownian motion
  - B. \*Non-commutatitve fractional Wishart process

# Part 0.A: Pioneering work of Eugene Wigner

# (Recalling Semicircle Law for Random Matrices)

#### 0.A. Random matrices and nuclear physics

Slow neutron resonance cross-sections on thorium 232 & uranium 238 nuclei. Energy(eV)



## 0.A. Gaussian Orthogonal Ensemble (GOE)

- Ensemble:  $\mathbf{Z} = (Z_n)$ ,  $Z_n$  is  $n \times n$  matrix with random entries.
- ► A) GOE: Z<sub>n</sub> = (Z<sub>n</sub>(j, k)) is n × n symmetric matrix with independent Gaussian entries in the upper triangular part:

$$Z_n(j,k) = Z_n(k,j) \sim N(0,1), \quad j \neq k,$$
  
$$Z_n(j,j) \sim N(0,2).$$

- ► B) Distribution of Z<sub>n</sub> is orthogonal invariant: Z<sub>n</sub>O<sup>T</sup> & Z<sub>n</sub> have same distribution for each orthogonal matrix O.
- Characterization GOE: A and B holds.

#### 0.A. Gaussian Orthogonal Ensemble (GOE)

• Joint density of eigenvalues of  $\lambda_1 > ... > \lambda_n$  of  $Z_n$ :

$$f_{\lambda_1,...,\lambda_n}(x_1,...,x_n) = k_n \underbrace{\left[\prod_{j=1}^n \exp\left(-\frac{1}{4}x_j^2\right)\right]}_{\text{independence}} \underbrace{\left[\prod_{j$$

Non-diagonal RM with density: eigenvalues are strongly dependent due to Vandermonde determinant:
 x = (x<sub>1</sub>, ..., x<sub>n</sub>) ∈ C<sup>n</sup>

$$\Delta(x) = \det\left(\left\{x_j^{k-1}\right\}_{j,k=1}^n\right) = \prod_{j < k} |x_j - x_k|.$$

#### 0.A. Wigner semicircle law

**Wigner 1950s**: Birth of RMT when both dimensions goes to  $\infty$ .

- A heavy nucleus is a liquid drop composed of many particles with unknown strong interactions,
- so a random matrix would be a possible model for the Hamiltonian of a heavy nucleus.
- Which random matrix should be used?
- $\lambda_1 > ... > \lambda_n$  eigenvalues of scaled GOE:  $X_n = Z_n / \sqrt{n}$ .
- Sample Spectral Distribution  $\widehat{F}_n^{X_n}$
- Limiting Spectral Distribution (LSD): *F*<sup>X<sub>n</sub></sup> goes, as n → ∞, to Semicircle distribution on (-2, 2)

$$w(x) = \frac{1}{2\pi}\sqrt{4-x^2}, \quad |x| \le 2.$$

Semicircle distribution is also called free Gaussian distribution

#### 0.A. Simulation of Wigner semicircle law



#### 0.A. Precise statement of Wigner semicircle law

#### Semicircle distribution approximates the spectral distribution

**Theorem**: For each continuous bounded function f and  $\epsilon > 0$ ,

$$\lim_{n\to\infty} \mathbb{P}\left(\left|\int f(x)d\widehat{F}_n^{X_n}(x) - \int f(x)w(dx)\right| > \varepsilon\right) = 0$$

where w(x) is the density of semicircle distribution on (-2, 2)

$$w(x) = \frac{1}{2\pi}\sqrt{4-x^2}, \quad |x| \le 2$$

- Good predictions for moderate dimension n
- Breakthrough work by Eugene Wigner: Ann. Math., 1955, 1957, 1958

## 0.A. Gaussian Unitary Ensemble (GUE)

Wigner law also holds

• GUE:  $Z_n = (Z_n(j, k))$  is  $n \times n$  Hermitian with independent Gaussian entries:

$$Z = \begin{pmatrix} Z_n(1,1) & Z_n(1,2) & \dots & Z_n(1,n) \\ \overline{Z}_n(1,2) & Z_n(2,2) & & \\ \overline{Z}_n(1,n) & & & Z_n(n,n) \end{pmatrix}$$

$$\operatorname{Re}\left(Z_n(j,k)\right) \sim \operatorname{Im}\left(Z_n(j,k)\right) \sim N(0,(1+\delta_{jk}))/2),$$

 $\operatorname{\mathsf{Re}}\left(Z_n(j,k)
ight)$  ,  $\operatorname{\mathsf{Im}}\left(Z_n(j,k)
ight)$  ,  $1\leq j\leq k\leq n$  ,

are independent random variables

Distribution of Z<sub>n</sub> is unitary invariant: Z<sub>n</sub>U\* & Z<sub>n</sub> have same distribution for each unitary non-random matrix U

#### 0.A. Universality

Wigner semicircle law holds for Wigner ensembles:

$$X_n(k,j) = X_n(j,k) = \frac{1}{\sqrt{n}} \begin{cases} Z_{j,k}, & \text{if } j < k \\ Y_j, & \text{if } j = k \end{cases}$$

 $\{Z_{j,k}\}_{j\leq k}$  ,  $\{Y_j\}_{j\geq 1}$  independent sequences of i.i.d. r.v. with

$$\mathbb{E}Z_{1,2}=\mathbb{E}Y_1=$$
0,  $\mathbb{E}Z_{1,2}^2=$ 1,  $\mathbb{E}Y_1^2<\infty$ 

- Whatever values the random entries take, the LSD (Semicircle) has bounded support
- Joint density of eigenvalues of a Wigner matrix is not easy

The Marchenko-Pastur Law
 B. For Sample Covariance Matrix
 Both, data dimension and sample size large

#### 0.B. Marchenko-Pastur law

Marchenko-Pastur (1967), Mat. Sb.

► 
$$H = H_{p \times n} = (Z_{j,k} : j = 1, ..., p, k = 1, ..., n)$$
 i.i.d. r.v.  
 $\mathbb{E}(Z_{1,1}) = 0, \mathbb{E}(|Z_{1,1}|^2) = 1, \mathbb{E}(|Z_{1,1}|^4) < \infty$ 

• Sample covariance matrix  $S_n = \frac{1}{n}HH^*$ , ESD  $\widehat{F}_p^{S_n} = \widehat{F}_p^{\frac{1}{n}HH^*}$ 

► If 
$$p/n \to c > 0$$
,  $\widehat{F}_p^{S_n}$  goes to MP distribution:  

$$\mu_c(dx) = \begin{cases} f_c(x)dx, & \text{if } 0 < c < 1\\ (1-c)\delta_0(dx) + f_c(x)dx, & \text{if } c \ge 1 \end{cases}$$

$$f_c(x) = rac{c}{2\pi x} \sqrt{(x-a)(b-x)} \mathbf{1}_{[a,b]}(x)$$
  
 $a = (1-\sqrt{c})^2, \ b = (1+\sqrt{c})^2$ 

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### 0.B. Simulation Marchenko-Pastur law



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0. Random matrices have been used in many fields

- Statistics
- Physics
- Number theory
- Biology
- Finances
- Engineering
- Computer vision
- Machine learning
- ▶ .....
- Prominently...together with free probability in wireless communications:

#### Time-varying random matrices: why?

Couillet & Debbah (2011), *Random Matrix Methods for Wireless Communications.* Chapter 19, Perspectives:

- Performance analysis of a typical network with users in motion according to some stochastic behavior, is not accessible to this date in the restrictive framework of random matrix theory.
- It is to be believed that random matrix theory for wireless communications may move on a more or less long-term basis towards random matrix process theory for wireless communications. Nonetheless, these random matrix processes are nothing new and have been the interest of several generations of mathematicians.

The work was initiated by Dyson in 1962

## Part I: From Matrix Bm to Free Bm

# (Time-varying random matrix models for the Free Brownian motion)

#### I. Hermitian Brownian motion ensemble

▶ 
$$\mathbf{B}(t) = (B_n(t))_{n \ge 1}, t \ge 0.$$

•  $B_n(t)$  is  $n \times n$  Hermitian Brownian motion:

$$B_n(t)=(b_{ij}(t)), t\geq 0,$$

$$\mathsf{Re}(b_{ij}(t)) \sim \mathsf{Im}(b_{ij}(t)) \sim \mathsf{N}(0, t(1+\delta_{ij}))/2,$$

where  $\text{Re}(b_{ij}(t))$ ,  $\text{Im}(b_{ij}(t))$ ,  $1 \le i \le j \le n$  are independent one-dimensional Brownian motions

►  $(\lambda_1(t), ..., \lambda_n(t))_{t \ge 0}$  process of eigenvalues of  $\{B_n(t)\}_{t \ge 0}$  $\lambda_1(t) \ge \lambda_2(t) \ge ... \ge \lambda_n(t)$ 

## I. Dyson-Brownian motion

Time dynamics of the eigenvalues, dimension n fixed

#### Dyson (1962):

a) If eigenvalues start at different positions, they never collide

$$\mathbb{P}\left(\lambda_1(t)>\lambda_2(t)>...>\lambda_n(t)\quad orall t>0
ight)=1$$

b) They satisfy the Stochastic Differential Equation (SDE)

$$\lambda_i(t) = \lambda_i(0) + W_i(t) + \sum_{j \neq i} \int_0^t rac{\mathrm{d}s}{\lambda_j(s) - \lambda_i(s)}, i = 1, ..., n$$

 $\forall t > 0$ , where  $W_1, ..., W_n$  are 1-dimensional independent Bms

- Brownian part + repulsion force (at any time t)
- Proof uses classical stochastic calculus and martingale techniques (Anderson, Guionnet, Zeitouni, 2010, Tao, 2012)

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## I. Key ideas for the proof

Where does the noncoliding force come from?

1. Eigenvalues are smooth functions of entries:

$$\lambda_i(t) = F(B_n(t)) = F((b_{jk}(t)))$$

2. Itô formula from classical stochastic calculus:  $b = (b_t)_{t \ge 0}$ 1-dimensional Bm

$$F(b_t) = \int_0^t F_x(b_s) \mathrm{d}b_s + \frac{1}{2} \int_0^t F_{xx}(b_s) \mathrm{d}s$$

3. Hadamard second variational formula: For a matrix A = A(x) depending smoothly on x, and  $Au_i = \lambda_i u_i$ ,  $u_i^* u_i = 1$ 

$$rac{d^2}{dx^2}\lambda_i = u_i^*A_{xx}u_i + 2\sum_{j
eq i}rac{1}{\lambda_i-\lambda_j}\left|u_j^*A_xu_j
ight|$$

Existence of independent Bms W<sub>1</sub>, ..., W<sub>n</sub> from Lévy's characterization of Bm

I. Time-varying Wigner theorem and law of Free Bm

Consider the Dyson spectral measure-valued processes

$$\mu_t^{(n)} = rac{1}{n} \sum_{j=1}^n \delta_{\left\{ \lambda_j(t) / \sqrt{n} 
ight\}}, \quad t \ge 0, \, n \ge 1$$

- ▶ Notation: For *f*  $\mu$ -integrable function  $\langle \mu, f \rangle = \int f(x)\mu(dx)$
- Uniform Wigner theorem

$$\mathbb{P}\left(\lim_{n\to\infty}\sup_{0\leq t\leq T}\left|\left\langle \mu_{t}^{(n)},f\right\rangle - \left\langle w_{t},f\right\rangle\right| = 0, \forall f\in C_{b}(\mathbb{R})\right) = 1$$

► The family of probability measures {w<sub>t</sub>}<sub>t≥0</sub> is the Law of the Free Brownian motion,

$$w_t(dx) = \frac{1}{2\pi t} \sqrt{4t - x^2} \mathbb{1}_{[-2\sqrt{t}, 2\sqrt{t}]}(x) dx$$

### I. Smooth vs non smooth SDE

Interacting systems

 SDE with both smooth drift & diffusion coefficients β and *α* are of the form

$$\begin{split} X_{n,i}(t) &= X_{n,i}(0) + \frac{1}{\sqrt{n}} \sum_{j \neq i} \int_0^t \beta(X_{n,j}(s), X_{n,i}(s)) dW_i^{(n)}(t) \\ &+ \frac{1}{n} \sum_{j \neq i} \int_0^t \alpha(X_{n,j}(s), X_{n,i}(s)) ds. \end{split}$$

While Dyson-Brownian motion has very non smooth drift

$$X_{n,i}(t) = X_{n,i}(0) + \frac{1}{\sqrt{n}} W_i^{(n)}(t) + \frac{1}{n} \sum_{j \neq i} \int_0^t \frac{1}{X_{n,i}(s) - X_{n,j}(s)} \mathrm{d}s$$

Empirical measure valued processes

$$\mu_t^{(n)} = rac{1}{n} \sum_{j=1}^n \delta_{X_{n,j}(t)}, \quad t \ge 0, n \ge 1$$

#### I. Smooth vs non smooth SDE

- 1. For Interacting SDE with both smooth drift & diffusion coefficients:
- McKean (Lect. Series Differ. Equat. 1967):  $\{\mu^{(n)}\}_{t\geq 0}$  converges weakly in probability to  $\{\mu_t\}_{t\geq 0}$ , which is the law of a Itô stochastic differential equation.
- 2. Interacting SDE with non-colliding forces arise from eigenvalue processes of matricial processes:
- [Bru (1989), Rogers & Shi (1993), Konig & O´Connell (2001), Cabanal-Duvillard & Guionnet (2001), Katori & Tanemura (2004)].
- ► The family of probabilities  $\{w_t, t \ge 0\}$  is not the law of a SDE equation, but the law of a noncommutative process: Free Brownian motion

## I. Fluctuations limit: Another difference

- 1. Interacting SDE with smooth drift & diffusion coefficients:
- Limits of fluctuations (CLT):  $S_t^{(n)} = \sqrt{n}(\mu_t^{(n)} \mu_t)$ .
- Hitsuda and Mitoma (JMA, 1986): S<sub>t</sub><sup>(n)</sup> converges to
   Gaussian process in nuclear space (Kallianpur & Pérez-Abreu (AMO, 1988), Kallianpur & Xiong (LNS, 1995))
- Gaussian fluctuations for interacting particle systems with a class of singular kernels, Wang, Zhao & Zhu (2021)
- 2. Interacting SDE with non-colliding eigenvalues:
- Due to noncoliding forces, need to consider fluctuations

$$Y_t^{(n)} = n\left(\mu_t^{(n)} - w_t\right)$$

- Israelson (SPA, 2001), Bender (SPA, 2008): Y<sub>t</sub><sup>(n)</sup> converges to Gaussian process in nuclear space
- Unterbergerg (SPA, 2018): Global fluctuations for 1D log-gas dynamics (generalized Dyson-Brownian motion)

## What is free Brownian motion?

What is the law of free Brownian motion?

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#### I. Notation

- Let  $\mathcal{P}(\mathbb{R})$  be the set of probability measures on  $\mathbb{R}$ .
- Let C (ℝ<sub>+</sub>, P(ℝ)) be the spaces of continuous functions from ℝ<sub>+</sub> → P(ℝ), with the topology of uniform convergence on compact intervals of ℝ<sub>+</sub>
- For µ ∈ P(ℝ) and a function f : ℝ → ℝ that is µ-integrable we write

$$\langle \mu, f \rangle = \int_{\mathbb{R}} f(x) \mu(\mathrm{d}x)$$

▶ Cauchy transform of  $\mu \in \mathcal{P}(\mathbb{R}) : G_{\mu}(z) : \mathbb{C}^+ \to \mathbb{C}^-$ 

$$\mathcal{G}_{\mu}(z) = \int_{\mathbb{R}} rac{1}{z-x} \mu(\mathrm{d} x), \quad z \in \mathbb{C}^+$$

I. Law of free Bm as limiting measure valued process

Theorem (Cabanal-Duvillard & Guionnet (2001)) If  $\mu_0^{(n)} \to \delta_0$ , the family  $(\mu_t^{(n)})_{t\geq 0}$  of measure valued-processes converges weakly in  $C(\mathbb{R}_+, \mathcal{P}(\mathbb{R}))$  to a unique continuous probability-measure valued function such that  $\forall f \in C_b^2(\mathbb{R})$ 

$$\langle \mu_t, f \rangle = f(0) + \frac{1}{2} \int_0^t \mathrm{d}s \int_{\mathbb{R}^2} \frac{f'(x) - f'(y)}{x - y} \mu_s(\mathrm{d}x) \mu_s(\mathrm{d}y).$$

Moreover,  $\mu_t = \mathrm{w}_t$ ,  $t \geq 0$ 

**Well known fact**: The family of laws  $(w_t)_{t\geq 0}$  is characterized by its Cauchy transforms  $(G_t)_{t\geq 0}$ , being the unique solution of

$$rac{\partial G_t(z)}{\partial t} = G_t(z) rac{\partial G_t(z)}{\partial z}, \quad t>0 \ G_0(z) = -rac{1}{z}, \quad z\in \mathbb{C}^+,$$

 $G_t(z) \in \mathbb{C}^+ \text{ for } z \in \mathbb{C}^+ \text{ \& } \lim_{\eta \to \infty} \eta \left| G_t(i\eta) \right| \leq \infty \forall t \geq 0 \text{ for } t \geq$ 

# Part II: Free Brownian motion and Non-commutatitve fractional Brownian motion

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II. Non-commutatitve probability spaces

A non-commutative probability space  $(\mathcal{A}, \varphi)$  is a unital algebra  $\mathcal{A}$  over  $\mathbb{C}$  with a linear functional  $\varphi : \mathcal{A} \to \mathbb{C}$  with  $\varphi(1_{\mathcal{A}}) = 1$ . Elements of  $\mathcal{A}$  are called *non-commutative random variables* 

#### Examples

1.  $\mathcal{A} = \mathbb{M}_d(\mathbb{C}) \ d imes d$  matrices with complex entries

$$\varphi(\cdot) = \operatorname{tr}_d(\cdot) = \frac{1}{d}\operatorname{tr}(\cdot)$$

2. 
$$\mathcal{A} = \mathcal{L}_{\infty}(\Omega, \mathcal{F}, \mathbb{P}),$$
  
 $\varphi(\cdot) = \mathbb{E}(\cdot)$   
3.  $\mathcal{A} = \mathbb{M}_{d}(\mathcal{L}_{\infty}(\Omega, \mathcal{F}, \mathbb{P})),$ 

$$\varphi(\cdot) = \mathbb{E}\mathrm{tr}_d(\cdot)$$

4.  $\mathcal{A} = L(\mathcal{H})$  algebra of linear operators on a Hilbert space,  $u \in \mathcal{H}, \|u\| = 1$  $\varphi(\cdot) = \langle \cdot u, u \rangle$ 

- We should think of φ as playing the role of the expectation in classical probability theory
- We talk about the moments of a, referring to the values of φ (a<sup>k</sup>), k ≥ 0
- More generally, for a tuple  $a_1, \ldots, a_n \in \mathcal{A}$ , the values

$$\varphi(\mathbf{a}_{i_1}^{m_1}...\mathbf{a}_{i_k}^{m_k})$$

for  $k \ge 0, 1 \le i_1, \ldots, i_k \le n, m_1, \ldots, m_k \ge 0$ , are known as the *joint moments* of  $a_1, \ldots, a_n$ 

When an algebraic distribution is given by an analytic distribution?

#### II. Non-commutatitve probability spaces

Generality needed to deal with free probability

**Remember classical case:** A real random variable R has distribution  $\mu$  on  $\mathbbm{R}$  iff

$$\mathbb{E}f(R) = \int_{\mathbb{R}} f(x)\mu(\mathrm{d}x), \quad \forall f \in B_b(\mathbb{R})$$

Non-commutatitve case needs:

(i) Given a p.m.  $\mu$  on  $\mathbb{R}$  with **bounded support**, there exist a  $C^*$ -probability space  $(\mathcal{A}, \varphi)$  and a self-adjoint  $\mathbf{a} \in \mathcal{A}$  with

$$\varphi(f(\mathbf{a})) = \int_{\mathbb{R}} f(x)\mu(\mathrm{d}x), \quad \forall f \in C_b(\mathbb{R})$$

(ii) **Given a p.m.**  $\mu$  on  $\mathbb{R}$ , there exists a  $W^*$ -probability space  $(\mathcal{A}, \varphi)$  and self-adjoint operator **a** on a Hilbert space H such that

$$f(\mathbf{a}) \in \mathcal{A} \quad \forall f \in B_b(\mathbb{R}),$$
 (1)

$$\varphi(f(\mathbf{a})) = \int_{\mathbb{R}} f(x)\mu(\mathrm{d}x), \quad \forall f \in B_b(\mathbb{R})$$

If (1) holds, it is said that **a** is affiliated with  $\mathcal{A}_{\downarrow}$ ,  $\mathcal{A}_{\downarrow}$ ,

## II. Free Random Variables

Definition

A family of subalgebras  $\{A_i\}_{i \in I} \subset A$  in a non-commutative probability space is *free* (freely independent) if

$$\varphi(\mathbf{a}_1\mathbf{a}_2\cdots\mathbf{a}_n)=\mathbf{0}$$

whenever  $\varphi(\mathbf{a}_j) = 0$ ,  $\mathbf{a}_j \in \mathcal{A}_{i_j}$ , and  $i_1 \neq i_2, i_2 \neq i_3, ..., i_{n-1} \neq i_n$ 

#### Definition

If  $\mathbf{a}_1, \mathbf{a}_2$  are freely independent, with distributions  $\mu_{a_1}$  and  $\mu_{a_2}$ , the distribution of  $\mathbf{a}_1 + \mathbf{a}_2$  is the free convolution  $\mu_{a_1} \boxplus \mu_{a_2}$ 

- ► Recall: If a<sub>1</sub>, a<sub>2</sub> are classical independent, the distribution of a<sub>1</sub> + a<sub>2</sub> is the classical convolution µ<sub>a1</sub> \* µ<sub>a2</sub>
- $\blacktriangleright$  If  $arphi = \mathbb{E}$  and there is commutativity

$$\mathbb{E}(\mathbf{a}_1\mathbf{a}_2\cdots\mathbf{a}_n)=\mathbb{E}(\mathbf{a}_1)\cdots\mathbb{E}(\mathbf{a}_n)=\mathbf{0}$$

II. Free independence allows to compute joint moments Example

Computation of  $\varphi(abab)$  when a & b are freely independent: Suppose  $\{a_1, a_3\}$  and  $\{a_2, a_4\}$  are freely independent. Since

$$arphi(\mathbf{a}_i-arphi(\mathbf{a}_i)\mathbf{1}_\mathcal{A})=\mathbf{0}_i$$

$$\begin{split} \varphi(\mathbf{a}_1 - \varphi(\mathbf{a}_1)\mathbf{1}_{\mathcal{A}})\varphi(\mathbf{a}_2 - \varphi(\mathbf{a}_2)\mathbf{1}_{\mathcal{A}})\varphi(\mathbf{a}_3 - \varphi(\mathbf{a}_3)\mathbf{1}_{\mathcal{A}})\varphi(\mathbf{a}_4 - \varphi(\mathbf{a}_4)\mathbf{1}_{\mathcal{A}}) = 0 \\ \text{Computations yield} \end{split}$$

$$\begin{split} \varphi(\mathbf{a}_1\mathbf{a}_2\mathbf{a}_3\mathbf{a}_4) &= \varphi(\mathbf{a}_1\mathbf{a}_3)\varphi(\mathbf{a}_2)\varphi(\mathbf{a}_4) + \varphi(\mathbf{a}_1)\varphi(\mathbf{a}_3)\varphi(\mathbf{a}_2\mathbf{a}_4) \\ &- \varphi(\mathbf{a}_1)\varphi(\mathbf{a}_2)\varphi(\mathbf{a}_3)\varphi(\mathbf{a}_4) \end{split}$$

In particular if  $\mathbf{a}_1 = \mathbf{a}_3 = \mathbf{a}$  and  $\mathbf{a}_2 = \mathbf{a}_4 = \mathbf{b}$ 

 $\varphi(\mathbf{a}\mathbf{b}\mathbf{a}\mathbf{b}) = \varphi(\mathbf{a})^2\varphi(\mathbf{b}^2) + \varphi(\mathbf{a}^2)\varphi(\mathbf{b})^2 - \varphi(\mathbf{a})^2\varphi(\mathbf{b})^2 \neq \varphi(\mathbf{a}^2)\varphi(\mathbf{b}^2)$ 

## II. Application: Free Central Limit Theorem

#### Theorem

Let  $\mathbf{a}_1, \mathbf{a}_2,...$  be a sequence of independent free random variables with the same distribution with all moments. Assume that  $\varphi(\mathbf{a}_1) = 0$  and  $\varphi(\mathbf{a}_1^2) = t$ . Then the distribution of

$$\mathbf{Z}_m = \frac{1}{\sqrt{m}}(\mathbf{a}_1 + \ldots + \mathbf{a}_m)$$

converges, as  $m \rightarrow \infty$ , to the semicircle distribution

$$\mathbf{w}_t(x) = rac{1}{2\pi}\sqrt{4t-x^2}, \quad |x| \leq 2\sqrt{t}$$

with moments  $m_{2k+1} = 0$  and  $m_{2k} = t^{2k} {\binom{2k}{k}}/{(k+1)}$ .

- Semicircle or Wigner distribution plays the role of classical Gaussian in free probability
- Marchenko-Pasture distribution (free Poisson) plays the role of Poisson distribution

#### II. RMT and Free Probability: Why useful?

From the Blog of Terence Tao (Free Probability):

- The significance of free probability to random matrix theory lies in the fundamental observation that random matrices which are independent in the classical sense, also tend to be independent in the free probabilistic sense, in the large limit.
- Because of this, many tedious computations in random matrix theory, particularly those of an algebraic or enumerative combinatorial nature, can be done more quickly and systematically by using the *framework of free probability*.
- Voiculescu (1991), Limit Laws for random matrices and free products. Invent. Math.

 Books on random matrices and wireless communications include free probability.

#### II. RMT and Free Probability: Why useful?

- ► Knowing eigenvalues of n × n random matrices X<sub>n</sub> & Y<sub>n</sub>, what are the eigenvalues of X<sub>n</sub> + Y<sub>n</sub>? X<sub>n</sub>Y<sub>n</sub>?
- ▶ In general if X<sub>n</sub> and Y<sub>n</sub> do not commute,

$$\lambda(X_n+Y_n)\neq\lambda(X_n)+\lambda(Y_n)$$

► However, if X<sub>n</sub> & Y<sub>n</sub> are asymptotically free, LSD of X<sub>n</sub> + Y<sub>n</sub> can be computed as (free convolution)

$$LSD(X_n + Y_n) = LSD(X_n) \boxplus LSD(Y_n)$$

The problem is similar to the computation to the distribution of the sum of two independent random variables: product of characteristic functions or moment generating functions (classical convolution)

## II. Free Brownian motion

A non-commutatitve process

A Free Brownian motion is a family  $S = \{S_t\}_{t \ge 0}$  of self-adjoint random variables in a non-commutative probability space  $(\mathcal{A}, \varphi)$  such that:

- 1.  $S_0 = 0$
- 2. For  $t_2 \geq t_1 \geq$  0,  $S_{t_2} S_{t_1}$  has law  $w_{t_2-t_1}$
- 3. For all  $n \ge 1$  and  $t_n > \cdots > t_1 > 0$ , the increments  $S_{t_n} S_{t_{n-1}}, \dots, S_{t_2} S_{t_1}, S_{t_1}$  are freely independent with respect to  $\varphi$
- For every t ≥ 0, St has semicircle law wt of zero mean and variance t > 0
- One has Stochastic calculus for the free Brownian motion (Anshelevich, JFA, 2002; Biane, FICAMS, 1997; Biane & Speicher, PTRF, 1998); Kemp, Nourdin, Peccati, Speicher, AP, 2012)

#### II. Semicircular process

Free Brownian motion is an example of a Semicircular process X = {X<sub>t</sub>}<sub>t≥0</sub> ⊂ A, self-adjoint random variables: For every k ≥ 1, t<sub>1</sub>, ..., t<sub>k</sub> ∈ [0, ∞) and θ<sub>1</sub>, ..., θ<sub>k</sub> ∈ ℝ, the non-commutative random variable θ<sub>1</sub>X<sub>t<sub>1</sub></sub> + · · · + θ<sub>k</sub>X<sub>t<sub>k</sub></sub> has Semicircle law

$$\mathbf{w}_{m,\sigma^2}(\mathbf{d}x) = \frac{1}{2\pi\sigma^2}\sqrt{4\sigma^2 - (x-m)^2}\mathbf{1}_{[m-2\sigma,m+2\sigma]}(x)\mathbf{d}x$$

for some  $m \in \mathbb{R}$ ,  $\sigma^2 > 0$ 

► The law of a centered semicircular processes (φ(X<sub>t</sub>) = 0 for every t > 0) is uniquely determined by its covariance function

$$\Gamma(s,t) = \varphi(X_t X_s)$$

Centered semicircular X = {X<sub>t</sub>}<sub>t≥0</sub> has stationary increments  $Γ(s, t) = Γ(|t - s|) = φ(X_{|t-s|})$ 

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#### II. Non-commutative Fractional Brownian Motion Nourdin and Taqqu (JTP, 2104)

Let H ∈ (0, 1). A noncommutative fractional Brownian motion (ncfBm) of Hurst parameter H is a centered semicircular process S<sup>H</sup> = {S<sup>H</sup><sub>t</sub>}<sub>t≥0</sub> in a non-commutative probability space (A, φ) with covariance function

$$\varphi(S_t^H S_s^H) = \frac{1}{2} \left( t^{2H} + s^{2H} - |t - s|^{2H} \right)$$

▶ For each t > 0,  $S_t^H$  has the semicircle law  $w_t^H$  on  $(-2t^H, 2t^H)$ 

$$w_t^H(dx) = rac{1}{2\pi t^{2H}} \sqrt{4t^{2H} - x^2} dx, \quad |x| \le 2t^H$$

ncfBm has stationary increments: For every t, s > 0

$$\varphi(\left(S_t^H - S_s^H\right)^2) = |t - s|^{2H}$$

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Part III: Extensions to other matrix and non-commutative processes

# III.A. From matrix fractional Bm to non-commutatitve fractional Bm

(Time-varying random matrix models for the non-commutatitve fractional Bm)

#### III.A. One-dimensional fractional Brownian motion

A one-dimensional fractional Brownian motion  $b^{H} = \{b^{H}(t)\}_{t \ge 0}$ is a zero-mean classical Gaussian process with covariance

$$\mathbb{E}(b^{H}(t)b^{H}(s))=rac{1}{2}\left(t^{2H}+s^{2H}-\left|t-s
ight|^{2H}
ight)$$

Stationary increments: For s, t > 0

$$\mathbb{E}\left|b^{H}(t)-b^{H}(s))\right|^{2}=\left|t-s\right|^{2H}$$

- ► Self-similarity:  $(a^{-H}b^{H}(at))_{t\geq 0} \stackrel{\text{law}}{=} (b^{H}(t))_{t\geq 0}$
- H = 1/2 is 1-dimensional Bm (independent increments)
- Itô stochastic calculus cannot be used for  $H \neq 1/2$
- Need classical fractional stochastic calculus: Skorohod, Young

#### III.A. Matrix fractional Brownian motion

Consider n(n+1)/2 independent 1-dimensional fractional Brownian motions with  $H \in (1/2, 1)$ 

$$\{\{b_{i,j}^{H}(t), t \ge 0\}, 1 \le i, j \le n\}$$

▶ n × n symmetric matrix fractional Brownian motion:

$$egin{aligned} \mathbf{B}_{n}^{H}(t) &= (B_{ij}^{H}(t))_{i,j=1}^{n} \ B_{ij}^{H}(t) &= b_{i,j}^{H} ext{ if } i < j \ B_{ii}^{H}(t) &= \sqrt{2} b_{i,i}^{H}(t) \end{aligned}$$

- ▶ For  $0 < t_1 < \cdots < t_p$ , the increments  $(\mathbf{B}_n^H(t_k t_{k-1}))_{n \ge 1}$ , k = 1, ..., p are not independent
- ▶ Let  $\lambda_1(t) \ge \lambda_2(t) \ge ... \ge \lambda_n(t)$  be the eigenvalues of  $\mathbf{B}_n^H(t)$

#### III.A. Matrix fractional Brownian motion

Nualart and Pérez-Abreu (SPA, 2014): no collision

$$\begin{split} H &> 1/2 \\ 1. \ \text{If } \lambda_1(0) > \lambda_2(0) > ... > \lambda_n(0)) \text{ the eigenvalues never collide:} \\ \mathbb{P} \left( \lambda_1(t) > \lambda_2(t) > ... > \lambda_n(t) \quad \forall t > 0 \right) = 1 \quad (*) \end{split}$$

2. For any t > 0 and i = 1, ..., n

$$egin{aligned} \lambda_i(t) &= \lambda_i(0) + Y_i(t) + 2H \sum_{j 
eq i} \int_0^t rac{1}{\lambda_i(s) - \lambda_j(s)} \mathrm{d}s \ Y_i(t) &= \sum_{k \leq h} \int_0^t rac{\partial \lambda_i(s)}{\partial b^H_{kh}(s)} \delta b^H_{kh}(s) \end{aligned}$$
 (\*\*)

- Stochastic integral in (\*\*) is in the sense of Skorohod. Classical Itô stochastic calculus cannot be used for H ≠ 1/2
- Proof of (\*) uses the Young stochastic integral
- Y<sub>i</sub>(t) is not a fractional Brownian motion, but it is a self-similar process: ∀ a > 0, (a<sup>-H</sup>Y<sub>i</sub>(at))<sub>t≥0</sub> <sup>law</sup> (Y<sub>i</sub>(t))<sub>t≥0</sub>

#### III.A. Matrix fractional Brownian motion

Jaramillo and Nualart (RMTA, 2021) and others: colliding case!!

Case H < 1/2

For T > 0

 $\mathbb{P}(\lambda_i(t) = \lambda_i(t) \text{ for some } t \in (0, T), \text{ and } 1 \leq i, j \leq n) = 1$ 

Tools: Hitting probabilities for Gaussian processes and capacity of Hausdorff dimension for measurable sets

General framework and some results of Jaramillo and Nualart:

- Matrix-valued Gaussian processes with a covariance R
- Fightharpoonup fBm real case: eigenvalues do not collide if H > 1/2 and collide if H < 1/2. Case H = 1/2 is Mckean (1969) method
- fBm complex case: eigenvalues do not collide if H > 1/3 and collide if H < 1/3
- Case H = 1/3, no collision: Lee, Song, Xiao and Yuan, 2023
- Collision of k eigenvalues: Song, Xiao and Yuan (2021)

## III.A. Time-varying Wigner theorem

Pardo, Pérez G, Pérez-Abreu (JTP, 2016)

Let H > 1/2. Consider the empirical spectral measure-valued processes of the re-scaled matrix fractional Bm  $\mathbf{B}_n^H(t)/\sqrt{n}$ 

$$\mu_t^{(n)} = \frac{1}{n} \sum_{j=1}^n \mathbf{1}_{\{\lambda_j(t)/\sqrt{n}\}}, \ t \ge 0, n \ge 1$$

1. Fix T > 0. For all continuous bounded function f and  $\varepsilon > 0$ 

$$\lim_{n \to \infty} \mathbb{P}\left(\sup_{0 \le t \le T} \left| \int f(x) d\mu_t^{(n)}(x) - \int f(x) w_t^H(x) dx \right| > \varepsilon \right) = 0$$

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where  $w_t^H$  is the semicircle distribution on  $(-2t^H, 2t^H)$ 

- 2. The family of measure-valued processes  $\{(\mu_t^{(n)})_{t\geq 0}: n\geq 1\}$  converges to  $(\mathbf{w}_t^H)_{t\geq 0}$ , the law of a non-commutative fractional Bm of Hurst parameter  $H\in (1/2, 1)$
- 3. Proof needs non collinding eigenvalues

#### III.A. Precise statement

Pardo, Pérez G, Pérez-Abreu (JTP, 2016): Functional Wigner

1. The family of measure-valued empirical spectral processes  $\{(\mu_t^{(n)})_{t\geq 0}: n\geq 1\}$  converges weakly in  $C(\mathbb{R}_+, \mathcal{P}(\mathbb{R}))$  to the unique continuous probability-measure valued function  $(\mu_t)_{t\geq 0}$  satisfying, for each  $t\geq 0$ ,  $f\in C_b^2(\mathbb{R})$ ,

$$\langle \mu_t, f \rangle = \langle \mu_0, f \rangle + H \int_0^t ds \int_{\mathbb{R}^2} \frac{f'(x) - f'(y)}{x - y} s^{2H-1} \mu_s(dx) \mu_s(dy)$$

Moreover  $\mu_t = \mathbf{w}_t^H$ 

2. The Cauchy transform  $G_t(z) = \int_{\mathbb{R}} \frac{\mu_t(dx)}{z-x}$  of  $\mu_t$  is the unique solution to the initial value problem

$$\left\{egin{aligned} &rac{\partial}{\partial t}G_t(z)=Ht^{2H-1}G_t(z)rac{\partial}{\partial z}G_t(z), & t>0,\ &G_0(z)=\int_{\mathbb{R}}rac{\mu_0(dx)}{z-x}, & z\in\mathbb{C}^+ \end{aligned}
ight.$$

**Extensions:** Matrices built from independent 1-dimensional SDE driven by fBm, H > 1/2 (Song, Yao, Yuan, 2020).

Jaramillo, Pardo, Pérez-G (JTP, 2019): Functional Wigner

Symmetric matrix built from independent Gaussian processes with regular covariance RIncludes fractional Brownian motion with index 0 < H < 1/2Proof does not need non-colliding eigenvalues and uses extended Skorohod integral

1.  $\{(\mu_t^{(n)})_{t\geq 0}: n\geq 1\}$  converges weakly in  $C(\mathbb{R}_+, \mathcal{P}(\mathbb{R}))$  to  $(\mu_t)_{t\geq 0}$  satisfying,  $t\geq 0$ ,  $f\in C_b^2(\mathbb{R})$ ,

$$\langle \mu_t, f \rangle = \langle \mu_0, f \rangle + \frac{1}{2} \int_0^t ds \int_{\mathbb{R}^2} \frac{f'(x) - f'(y)}{x - y} R'(s, s) \mu_s(dx) \mu_s(dy)$$

2. The Cauchy transform  $G_t$  of  $\mu_t$  satisfies  $G_t(z) = F_{R(t,t)}(z)$ ,  $t > 0, z \in \mathbb{C}^+$ , with F the solution of

$$\frac{d}{d\tau}F_t(z)=F_t(z)\frac{d}{dz}F_t(z),$$

 $F_0(z)$  the Cauchy transform of  $\mu_0, t > 0, z \in \mathbb{C}^+$ 

Diaz, Jaramillo, Pardo (AIHPB, 2022): Gaussian fluctuations

Symmetric matrix from independent Gaussian processes with regular covariance  ${\cal R}$ 

$$Z_{\Phi}^{(n)}(t) = \left( (Z_{\phi_1}^{(n)}(t), ..., Z_{\phi_r}^{(n)}(t)); t \ge 0 \right)$$

$$\begin{split} \Phi &= (\phi_1, ..., \phi_r), \ r \geqslant 1, \phi_i \text{ in a large class of test functions} \\ Z_{\phi}^{(n)}(t) &= n \left[ \int_{\mathbb{R}} \phi(t) \mu_t^{(n)}(dx) - \mathbb{E}\left( \int_{\mathbb{R}} \phi(t) \mu_t^{(n)}(dx) \right) \right] \end{split}$$

- 1. There exists a continuous  $\mathbb{R}^r$ -valued centered Gaussian process  $\Lambda_{\Phi}(t) = ((\Lambda_{\phi_1}(t), ..., \Lambda_{\phi_r}(t)); t \ge 0)$  such that  $Z_{\phi}^{(n)}(t)$  converges stably to  $\Lambda_{\Phi}(t)$ , in the topology of uniform convergence over compact sets
- 2. An explicit expression for covariance of  $\Lambda_{\Phi}(t)$  is found
- 3. An upper bound for the total variation distance between the laws of  $Z_{\Phi}{}^{(n)}(t)$  and  $\Lambda_{\Phi}(t)$  is established

Diaz, Jaramillo, Pardo (AIHPB, 2022): Gaussian fluctuations

#### **Proof uses**

 CLT in the Wiener chaos (related to Fourth Moment Problem of Nualart & Peccati (Ann Probab, 2005))

Diaz, Jaramillo, Pardo (AIHPB, 2022): Gaussian fluctuations

#### **Proof uses**

 CLT in the Wiener chaos (related to Fourth Moment Problem of Nualart & Peccati (Ann Probab, 2005))

Free independence and multiple Wigner integrals

Diaz, Jaramillo, Pardo (AIHPB, 2022): Gaussian fluctuations

#### **Proof uses**

 CLT in the Wiener chaos (related to Fourth Moment Problem of Nualart & Peccati (Ann Probab, 2005))

- Free independence and multiple Wigner integrals
- Weak convergence results

Diaz, Jaramillo, Pardo (AIHPB, 2022): Gaussian fluctuations

#### **Proof uses**

 CLT in the Wiener chaos (related to Fourth Moment Problem of Nualart & Peccati (Ann Probab, 2005))

- Free independence and multiple Wigner integrals
- Weak convergence results
- Stein methods

Diaz, Jaramillo, Pardo (AIHPB, 2022): Gaussian fluctuations

#### **Proof uses**

 CLT in the Wiener chaos (related to Fourth Moment Problem of Nualart & Peccati (Ann Probab, 2005))

- Free independence and multiple Wigner integrals
- Weak convergence results
- Stein methods
- Malliavan Calculus

Part III.B: From Fractional Wishart process to Non-commutatitve Wishart process

Extensions of free Wishart process

#### III.B. Fractional Wishart process

•  $m, n \ge 1, m \times n$  matrix process

$$\{B_{m,n}(t)\}_{t\geq 0} = \left\{ \left( b_{m,n}^{j,k}(t) \right)_{1\leq j\leq m, 1\leq k\leq n} \right\}_{t\geq 0},$$

 $\left\{ \operatorname{Re}\left(b_{m,n}^{j,k}\left(t\right)\right) \right\}_{t\geq 0} \& \left\{ \operatorname{Im}\left(b_{m,n}^{j,k}\left(t\right)\right) \right\}_{t\geq 0} \text{ independent}$  1-dimensional fractional Bm of parameter  $H \in [1/2, 1)$ .

 Fractional Laguerre, fractional Wishart process: n × n matrix-valued process

$$L_{m,n}(t) = B^*_{m,n}(t)B_{m,n}(t), t \ge 0$$

▶  $0 \le \lambda_n(t) \le \cdots \le \lambda_1(t)$  eigenvalues of  $L_{m,n}(t)/n$ 

For  $H \in [1/2, 1)$  the noncoliding property holds

$$\mathbb{P}\left(\lambda_1(t) > \lambda_2(t) > ... > \lambda_n(t) > 0 \quad \forall t > 0\right) = 1$$

#### III.B. Fractional Wishart process

- ► *H* = 1/2:
  - Bru (JMA, 1989): noncoliding property and stochastic dynamics

$$egin{aligned} &\mathrm{d}\lambda_i(t) = \lambda_i(0) + rac{1}{\sqrt{n}}\sqrt{2\lambda_i(t)}\,W_i(t) \ &+ rac{1}{n}\int_0^t \left(m + \sum_{j
eq i}rac{\lambda_i(s) + \lambda_j(s)}{\lambda_i(s) - \lambda_j(s)}
ight)\mathrm{d}s,\ 1\leq i\leq n \end{aligned}$$

- ▶ Cabanal-Duvillard & Guionnet (AP, 2001), Pérez-Abreu & Tudor (EJP, 2009): limiting measure-valued process, when  $n/m \rightarrow c > 0$ , is dilation of free Poisson law
- ▶  $H \in (1/2, 1)$ : Pardo, Pérez G., Pérez-Abreu (JFA, 2017):
  - Noncoliding, stochastic dynamics of eigenvalues
  - Limiting measure valued process is fractional dilation of free Poisson law

#### III.B. Dilation rather than the law of free Poisson Law of non-commutative fractional Wishart process

The limit, when n/m o c > 0, of  $\mu_t^{(n)} = rac{1}{n} \sum_{j=1}^n \delta_{\lambda_j(t)}, \, t \geq 0$  ,

▶ is not the law of a Free Poisson process {m<sub>ct,b</sub>}<sub>t>0</sub>,

$$m_{a,b}(dx) = \begin{cases} f_{a,b}(x)dx, & a \ge 1\\ (1-a)\delta_0(dx) + f_{a,b}(x)dx, & 0 \le a < 1, \end{cases}$$

$$f_{a,b}(x) = \frac{1}{2\pi bx} \sqrt{4ab^2 - (x - b(1 + a))^2} \mathbf{1}_{[b(1 - \sqrt{a})^2, b(1 + \sqrt{a})^2]}(x)$$

rather fractional dilations of m<sub>c,1</sub>: µ<sup>H</sup><sub>c</sub>(t) = m<sub>c,1</sub> ∘ h<sup>-1</sup><sub>t</sub>, for h<sub>t</sub>(x) = t<sup>2H</sup>x, i.e.
µ<sup>H</sup><sub>c</sub>(t) = m<sub>c,t<sup>2H</sup></sub>

#### III.B. Characterization of the law

Cabanal-Duvillard & Guionnet (AP, 2001): H = 1/2Pardo, Pérez G, Pérez-Abreu (JFA, 2017):  $H \in (1/2, 1)$ 

#### Theorem

The family  $(\mu_c^H(t), t \ge 0)$  is characterized by the property that its Cauchy transform  $G_{c,H}$  is the unique solution to

$$\begin{aligned} \frac{\partial G_{c,H}}{\partial t}(t,z) &= 2Ht^{2H-1} \left[ \begin{array}{c} G_{c,H}^2(t,z) + \\ (1-c+2zG_{c,1/2}(t,z))\frac{\partial G_{c,H}}{\partial z}(t,z) \end{array} \right], t > 0\\ G_{c,H}(0,z) &= \int_{\mathbb{R}} \frac{\mu_{c,H}(0)(dx)}{x-z} \end{aligned}$$

**Extensions:** Wishart matrices built from independent 1-dimensional SDE driven by fBm, H > 1/2 (Song, Yao, Yuan, 2021).

Dyson-Brownian motion and other noncoliding processes

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# Thanks